

Nonlinear Systems

Lecture 11

02/26/13

Last time:

- Lyapunov $\left\{ \begin{array}{l} \text{stability} \\ \text{Functions (i.e. indirect method)} \end{array} \right.$

main result: $\dot{x} = f(x); f(0) = 0$

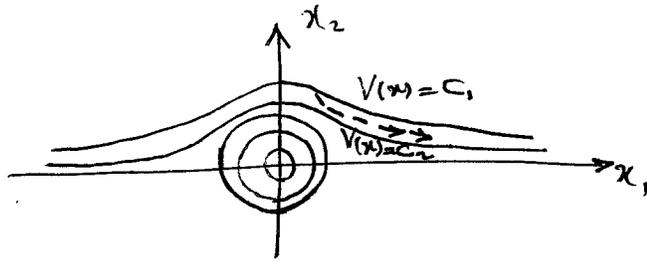
$V: D \rightarrow \mathbb{R}$ positive definite
($V(0) = 0, V(x) > 0$ for all $x \in D$)

1. $\dot{V}(x) \leq 0$, for all $x \in D \Rightarrow \bar{x} = 0$ is stable
(negative semidefinite)

2. $\dot{V}(x) < 0$ for all $x \in D \setminus \{0\} \Rightarrow \bar{x} = 0$ is LAS
 $V(0) = 0$

GAS: $\left. \begin{array}{l} V: \text{globally positive definite} \\ \text{radially unbounded} \\ (\|x\| \rightarrow \infty \Rightarrow V(x) \rightarrow \infty) \\ \oplus \\ \dot{V}: \text{globally negative definite} \end{array} \right\} \Omega_c = \{x; V(x) < c\}$

$$V(x) = \frac{x_1^2}{x_1^2 + 1} + x_2^2 \quad c_1 < c_2$$

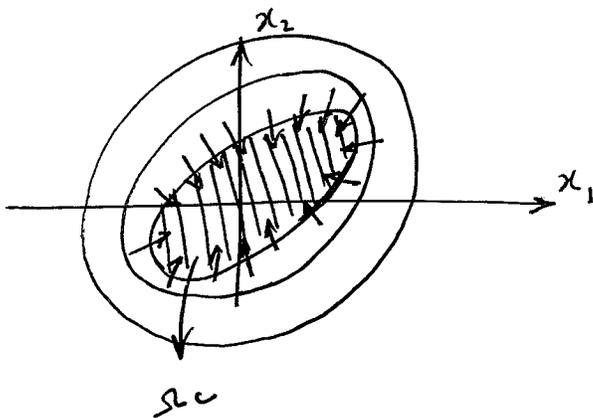


this example shows why radial unboundedness is crucial for Global AS.

1. Stability follows from positive invariance of Ω_c

$$\dot{V}(x) = \frac{\partial V}{\partial x} \dot{x} = \nabla V(x) \cdot f(x) \leq 0$$

$$x_0 \in \Omega_c \Rightarrow x(t) \in \Omega_c \text{ for all } t$$

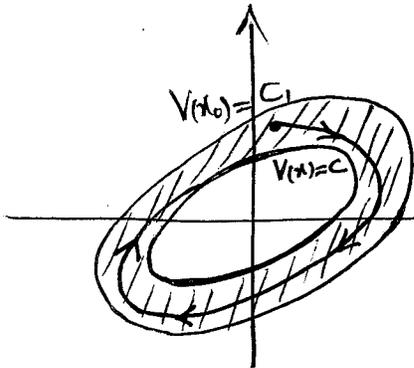


2. (Sketch of the proof) $\dot{V}(x) < 0, \forall x \in D \setminus \{0\}$

$\frac{dV}{dt} < 0 \Rightarrow V(t)$ is decreasing function bounded from below by zero

$$\Rightarrow \exists c \text{ st. } \lim_{t \rightarrow \infty} V(t) = c$$

assume $c > 0$ and show that this is not possible.



$$\max \dot{V} = -\gamma < 0$$

$$\{x; c < V(x) < C_1\}$$

$$V(x(t)) = V(x_0) + \int_0^t \underbrace{\dot{V}(x(\tau))}_{-\gamma} d\tau \quad \downarrow$$

$$\leq V(x_0) - \gamma t$$

\Rightarrow For t large enough $V(x(t))$ will become negative

\downarrow

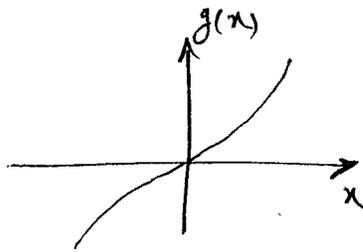
contradiction

Thus c has to be zero.

Examples

$$1) \quad \dot{x} = -g(x) \quad ; \quad x(t) \in \mathbb{R}$$

1a)



$$g(0) = 0$$

$$x \cdot g(x) > 0, \forall x$$

e.g.

$$g(x) = \underset{k}{\sqrt{x}} ; \Rightarrow g(0) = 0$$

$$x \cdot g(x) = \underset{k}{x^2} > 0 \quad \forall x \in \mathbb{R} \setminus \{0\}$$

$k > 0$

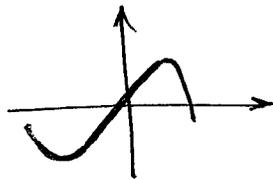
e.g.

$$g(x) = x^{2n+1} \quad n \in \mathbb{N} = \{1, 2, 3, \dots\}$$

$$\Rightarrow g(0) = 0$$

$$x \cdot g(x) = x^2 \cdot x^{2n} > 0 \quad \forall x \in \mathbb{R} \setminus \{0\}$$

1b)



$$g(0) = 0$$

$$x \cdot g(x) > 0 \quad \forall x \in (-b, c) \setminus \{0\}$$

or $(-a, a)$

e.g.

$$g(x) = \sin(x) ; b = c = \pi$$

$$g(x) = x - x^3 ; b = c = 1$$

* How to choose Lyapunov functions?

start with

$$V(x) = \frac{1}{2} x^2$$

→ g.p.d. (globally positive definite)
→ radially unbounded

$$\dot{V} = \frac{1}{2} 2x \cdot \dot{x} = x(-g(x)) = -xg(x) < 0$$

$$\frac{\partial V}{\partial x} \quad \downarrow \quad f(x)$$

for all $x \in D \setminus \{0\}$

$$D = \begin{cases} \mathbb{R} \text{ in (a)} \Rightarrow \text{GAS} & (\text{Global}) \\ (-b, c) \text{ in (b)} \Rightarrow \text{LAS} & (\text{Local}) \end{cases}$$

Alternative choice:

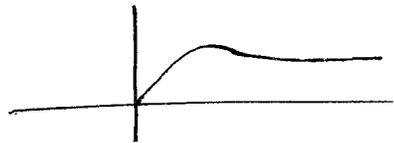
$$V(x) = \int_0^x g(\xi) d\xi$$

(Note! if $g(x) = x \Rightarrow V(x) = x^2$)

$$\dot{V}(x) = \frac{\partial V}{\partial x} f(x) = g(x) \cdot (-g(x)) = -g^2(x)$$

for 1(a) as long as we have

we will have radial unboundedness



$$\forall x \in \begin{cases} \mathbb{R} \setminus \{0\} \text{ in (a)} \Rightarrow \text{GAS} \\ (-b, c) \text{ in (b)} \Rightarrow \text{LAS} \end{cases}$$

$$2) \quad \dot{x}_1 = x_2$$

$$\dot{x}_2 = -ag(x) - bx_2$$

Pendulum special case of this with $g(x) = \sin x$

$$\text{Last time: } V(x) = a \int_0^{x_1} g(\xi) d\xi + \frac{1}{2} x_2^2$$

$$\dot{V} = \left[\frac{\partial V}{\partial x_1} \quad \frac{\partial V}{\partial x_2} \right] \begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \end{bmatrix} = \begin{bmatrix} ag(x_1) & x_2 \end{bmatrix} \begin{bmatrix} x_2 \\ -ag(x_1) - bx_2 \end{bmatrix}$$

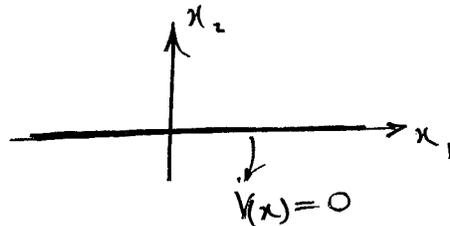
$$= \cancel{ax_2} g(x_1) - \cancel{ax_2} g(x_1) - bx_2^2$$

$$\Rightarrow \dot{V}(x) = -0x_1^2 - bx_2^2$$

$b = 0 \Rightarrow$ no damping (energy conserved)

$b > 0 \Rightarrow \dot{V}$ is negative semi-definite

choose $x_2 = 0$



$$\dot{V}(x) \leq 0 \text{ for all } x \in D \setminus \{0\}$$

\Downarrow

From our main result \Rightarrow ^{we can} \dot{V} only conclude stability in the sense of Lyapunov, not asymptotic stability.

~~///~~ We know that pendulum would eventually settle down (if $b > 0$)

\hookrightarrow there is viscous damping

Q: How can we sort out this issue?

Two options:

- choose different Lyapunov function (later)
- LaSalle's invariance principle

LaSalle's Invariance Principle

- applicable to time invariant systems
- can conclude asymptotic stability even when $\dot{V}(x) \leq 0$

Let $\Omega_c = \{x; V(x) \leq c\}$ be bounded and let $\dot{V}(x) \leq 0$ in Ω_c .

Define $S := \{x \in \Omega_c; \dot{V}(x) = 0\}$ and let M be

the largest invariant set in S . Then for every

$x(0) \in \Omega_c \Rightarrow x(t) \rightarrow M$. \blacksquare

In our example

$$V(x) = -bx_2^2 = 0 \Rightarrow x_2 = 0$$

$$S = \{x_1 \in (-b, c), x_2 = 0\}$$

$$x_2 \equiv 0 \Rightarrow \dot{x}_2 \equiv 0$$

↓

$$ag(x_1) - b \underset{0}{x_2} \equiv 0$$

$$g(x_1) \equiv 0 \Rightarrow (\text{can only happen for } x_1 = 0)$$

