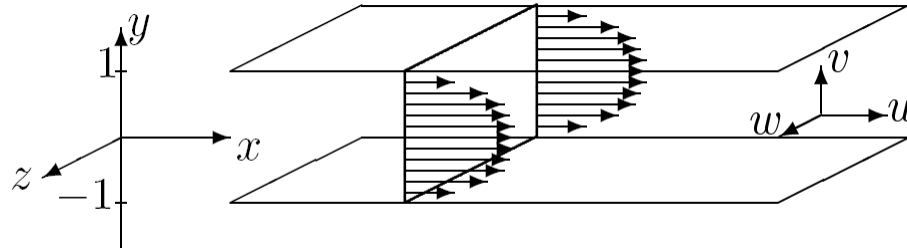


Lectures 13 & 14: ... and a bit of fluids

- Themes:
 - ★ Linearized Navier-Stokes (NS) equations in a channel flow
 - ★ Inner product that induces kinetic energy
 - ★ Non-normal nature of the dynamical generator
 - ★ Riesz spectral basis
- Approach: informal discussion using tools that we've learned so far (more later in the course)

Channel flow



- Steady-state solution: $[U(y) \ 0 \ 0]^T$
- Linearized NS and continuity equations

$$u_t + U(y) u_x + U'(y) v = -p_x + \frac{1}{Re} \Delta u$$

$$v_t + U(y) v_x = -p_y + \frac{1}{Re} \Delta v$$

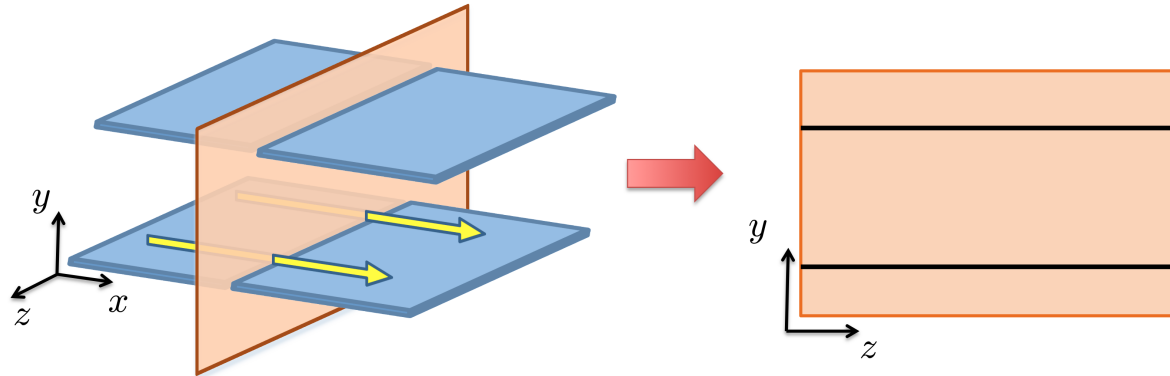
$$w_t + U(y) w_x = -p_z + \frac{1}{Re} \Delta w$$

$$u_x + v_y + w_z = 0$$

$$U(y) = \begin{cases} 1 - y^2, & \text{pressure driven flow} \\ y, & \text{shear driven flow} \end{cases}$$

$$U'(y) = \frac{dU(y)}{dy} \quad \Delta = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

Streamwise constant fluctuations



- Set $(\cdot)_x = 0$

$$u_t = -U'(y)v + \frac{1}{Re} \Delta u$$

$$v_t = -p_y + \frac{1}{Re} \Delta v$$

$$w_t = -p_z + \frac{1}{Re} \Delta w$$

$$0 = v_y + w_z$$

- ★ Define: stream-function in the (y, z) -plane

$$\{v = \psi_z, w = -\psi_y\}$$

- ★ Eliminate pressure from the equations

- ★ Rewrite equations in terms of

$$\phi = \begin{bmatrix} \psi & u \end{bmatrix}^T$$

Evolution model

$$\begin{bmatrix} \psi_t(t) \\ u_t(t) \end{bmatrix} = \begin{bmatrix} (1/Re) \mathcal{L} & 0 \\ \mathcal{C}_p & (1/Re) \mathcal{S} \end{bmatrix} \begin{bmatrix} \psi(t) \\ u(t) \end{bmatrix}$$

Orr-Sommerfeld: $\mathcal{L} = \Delta^{-1} \Delta^2$

Squire: $\mathcal{S} = \Delta$

Coupling: $\mathcal{C}_p = -U'(y) \partial_z$

- After Fourier transform in z

Laplacian: $\Delta = \partial_{yy} - k_z^2$

"Square of Laplacian": $\Delta^2 = \partial_{yyyy} - 2k_z^2 \partial_{yy} + k_z^4$

Coupling: $\mathcal{C}_p = -jk_z U'(y)$

Boundary conditions:

★ Dirichlet: $u(y = \pm 1, k_z, t) = 0$

★ Dirichlet and Neumann: $\psi(y = \pm 1, k_z, t) = \psi_y(y = \pm 1, k_z, t) = 0$

- Re-scale time: $\tau = t/Re$

$$\begin{bmatrix} \psi_\tau(\tau) \\ u_\tau(\tau) \end{bmatrix} = \begin{bmatrix} \mathcal{L} & 0 \\ Re \mathcal{C}_p & \mathcal{S} \end{bmatrix} \begin{bmatrix} \psi(\tau) \\ u(\tau) \end{bmatrix}$$

Inner product:

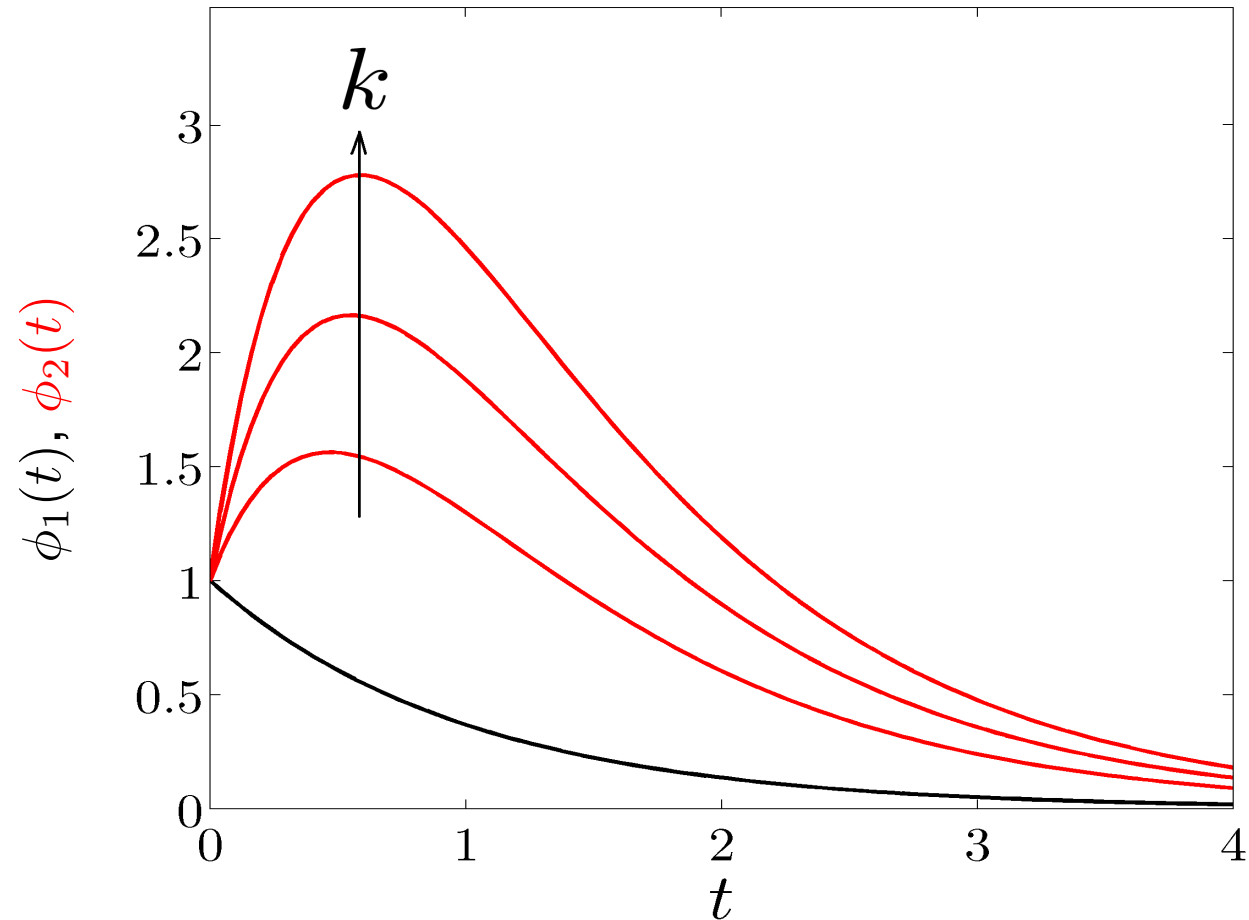
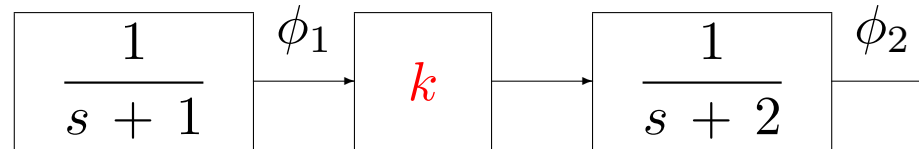
$$\begin{aligned} \langle \phi_1, \phi_2 \rangle_e &= \left\langle \begin{bmatrix} \psi_1 \\ u_1 \end{bmatrix}, \begin{bmatrix} \psi_2 \\ u_2 \end{bmatrix} \right\rangle_e \\ &= \left\langle \begin{bmatrix} \psi_1 \\ u_1 \end{bmatrix}, \begin{bmatrix} -\Delta & 0 \\ 0 & I \end{bmatrix} \begin{bmatrix} \psi_2 \\ u_2 \end{bmatrix} \right\rangle \\ &= \langle \psi_1, -\Delta \psi_2 \rangle + \langle u_1, u_2 \rangle \end{aligned}$$

Energy:

$$\begin{aligned} E &= \frac{1}{2} (\langle u, u \rangle + \langle v, v \rangle + \langle w, w \rangle) \\ &= \frac{1}{2} (\langle u, u \rangle + \langle \psi, -\Delta \psi \rangle) \end{aligned}$$

A finite dimensional example

$$\begin{bmatrix} \dot{\phi}_1 \\ \dot{\phi}_2 \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ k & -2 \end{bmatrix} \begin{bmatrix} \phi_1 \\ \phi_2 \end{bmatrix}$$



$$\dot{\phi}(t) = A \phi(t), \quad A A^* \neq A^* A$$

Let A have a full set of linearly independent e-vectors

$$A v_i = \lambda_i v_i \quad \Leftrightarrow \quad A \underbrace{\begin{bmatrix} v_1 & \cdots & v_n \end{bmatrix}}_V = \underbrace{\begin{bmatrix} v_1 & \cdots & v_n \end{bmatrix}}_V \underbrace{\begin{bmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_n \end{bmatrix}}_\Lambda$$

$$A^* w_i = \bar{\lambda}_i w_i \quad \Leftrightarrow \quad A^* \underbrace{\begin{bmatrix} w_1 & \cdots & w_n \end{bmatrix}}_W = \underbrace{\begin{bmatrix} w_1 & \cdots & w_n \end{bmatrix}}_W \underbrace{\begin{bmatrix} \bar{\lambda}_1 & & \\ & \ddots & \\ & & \bar{\lambda}_n \end{bmatrix}}_{\bar{\Lambda}}$$

choose w_i such that $w_i^* v_j = \delta_{ij}$



• A – diagonalizable:

$$A = \begin{bmatrix} v_1 & \cdots & v_n \end{bmatrix} \begin{bmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_n \end{bmatrix} \begin{bmatrix} w_1^* \\ \vdots \\ w_n^* \end{bmatrix}$$

- Action of A on $f \in \mathbb{C}^n$

$$\begin{aligned}
 Af &= \begin{bmatrix} v_1 & \cdots & v_n \end{bmatrix} \begin{bmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_n \end{bmatrix} \begin{bmatrix} w_1^* \\ \vdots \\ w_n^* \end{bmatrix} f \\
 &= \begin{bmatrix} v_1 & \cdots & v_n \end{bmatrix} \begin{bmatrix} \lambda_1 w_1^* f \\ \vdots \\ \lambda_n w_n^* f \end{bmatrix} \\
 &= \lambda_1 v_1 w_1^* f + \cdots + \lambda_n v_n w_n^* f \\
 &= \sum_{i=1}^n \lambda_i v_i \langle w_i, f \rangle
 \end{aligned}$$

- Solution to $\dot{\phi}(t) = A\phi(t)$

$$\phi(t) = e^{At} \phi(0) = \sum_{i=1}^n e^{\lambda_i t} v_i \langle w_i, \phi(0) \rangle$$

- E-value decomposition of $A = \begin{bmatrix} -1 & 0 \\ k & -2 \end{bmatrix}$

$$\{\lambda_1 = -1, \lambda_2 = -2\} \quad \left\{ v_1 = \frac{1}{\sqrt{1+k^2}} \begin{bmatrix} 1 \\ k \end{bmatrix}, v_2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix} \right\}$$

$$\left\{ w_1 = \begin{bmatrix} \sqrt{1+k^2} \\ 0 \end{bmatrix}, w_2 = \begin{bmatrix} -k \\ 1 \end{bmatrix} \right\}$$

- Solution to $\dot{\phi}(t) = A\phi(t)$

$$\begin{aligned} \phi(t) &= (e^{-t} v_1 w_1^* + e^{-2t} v_2 w_2^*) \phi(0) \\ &= \begin{bmatrix} e^{-t} & 0 \\ k(e^{-t} - e^{-2t}) & e^{-2t} \end{bmatrix} \begin{bmatrix} \phi_1(0) \\ \phi_2(0) \end{bmatrix} \end{aligned}$$

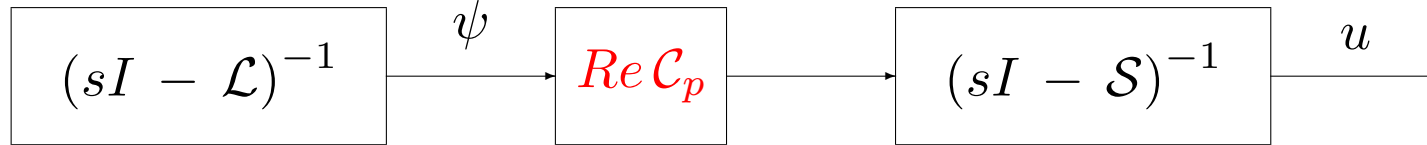
Back to fluids

$$\begin{bmatrix} \psi_\tau(\tau) \\ u_\tau(\tau) \end{bmatrix} = \begin{bmatrix} \mathcal{L} & 0 \\ \text{Re } \mathcal{C}_p & \mathcal{S} \end{bmatrix} \begin{bmatrix} \psi(\tau) \\ u(\tau) \end{bmatrix}$$

Orr-Sommerfeld

Coupling

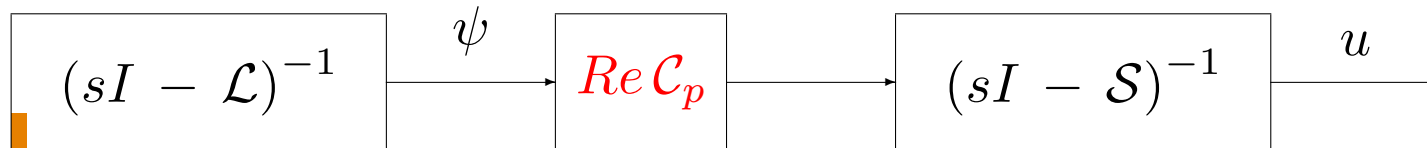
Squire



'glorified diffusion'

vortex tilting

viscous dissipation



- Adjoint of \mathcal{A} (w.r.t. $\langle \cdot, \cdot \rangle_e$):

$$\mathcal{A} = \begin{bmatrix} \mathcal{L} & 0 \\ \text{Re } \mathcal{C}_p & \mathcal{S} \end{bmatrix} \Rightarrow \left\{ \mathcal{A}^\dagger = \begin{bmatrix} \mathcal{L} & \text{Re } \mathcal{C}_p^\dagger \\ 0 & \mathcal{S} \end{bmatrix}, \mathcal{C}_p^\dagger = -jk_z \Delta^{-1} U'(y) \right\}$$

☞ \mathcal{A} : not normal \Leftrightarrow not diagonalizable by a unitary coordinate transformation

Spectral decomposition of \mathcal{A} and \mathcal{A}^\dagger

$$\begin{bmatrix} \mathcal{L} & 0 \\ \text{Re } \mathcal{C}_p & \mathcal{S} \end{bmatrix} \begin{bmatrix} \psi \\ u \end{bmatrix} = \lambda \begin{bmatrix} \psi \\ u \end{bmatrix} \Rightarrow \begin{cases} \mathcal{L} \psi = \lambda \psi \\ \mathcal{S} u = \lambda u - \text{Re } \mathcal{C}_p \psi \end{cases}$$

- Two sets of eigenvalues

$$(\lambda I - \mathcal{L}) \text{ not one-to-one} \Rightarrow \left\{ \lambda_{os}, \begin{bmatrix} \psi_{os} \\ u_{os} \end{bmatrix} \right\}$$

$$(\lambda I - \mathcal{S}) \text{ not one-to-one} \Rightarrow \left\{ \lambda_{sq}, \begin{bmatrix} 0 \\ u_{sq} \end{bmatrix} \right\}$$

- Homework:

★ fill in details for the e-value decomposition of \mathcal{A} and \mathcal{A}^\dagger

$$\text{Orr-Sommerfeld: } \begin{cases} \mathcal{L} \psi_{os} = \lambda_{os} \psi_{os}, & \psi_{os}(\pm 1) = \psi'_{os}(\pm 1) = 0 \\ \mathcal{S} u_{os} = \lambda_{os} u_{os} - \text{Re } \mathcal{C}_p \psi_{os}, & u_{os}(\pm 1) = 0 \end{cases}$$

$$\text{Squire: } \left\{ \lambda_{sq} = - \left(\left(\frac{n\pi}{2} \right)^2 + k_z^2 \right), \begin{bmatrix} 0 \\ u_{sq} \end{bmatrix} = \begin{bmatrix} 0 \\ \sin \left(\frac{n\pi}{2} (y + 1) \right) \end{bmatrix} \right\}$$

★ show that \mathcal{A} is a Riesz-spectral operator

Riesz-spectral operator

- Action of \mathcal{A} on $f \in \mathbb{H}$

$$[\mathcal{A}f](y) = \sum_{n=1}^{\infty} \lambda_{os,n} v_{os,n}(y) \langle w_{os,n}, f \rangle_e + \sum_{n=1}^{\infty} \lambda_{sq,n} v_{sq,n}(y) \langle w_{sq,n}, f \rangle_e$$

- Solution to $\phi_\tau(\tau) = \mathcal{A}\phi(\tau)$, $\phi(0) = f$

$$\phi(y, \tau) = \sum_{n=1}^{\infty} e^{\lambda_{os,n} \tau} v_{os,n}(y) \langle w_{os,n}, f \rangle_e + \sum_{n=1}^{\infty} e^{\lambda_{sq,n} \tau} v_{sq,n}(y) \langle w_{sq,n}, f \rangle_e$$

- Dependence of $u(y, k_z, \tau)$ on $\psi(y, k_z, 0) = \sum_{n=1}^{\infty} \alpha_n(k_z) \psi_{os,n}(y, k_z)$

$$u(y, k_z, \tau) = \operatorname{Re} \sum_{n=1}^{\infty} \left(\alpha_n e^{\lambda_{os,n} \tau} u_{os,n}(y, k_z, \tau) - \sum_{m=1}^{\infty} \frac{\alpha_m}{\lambda_{os,m} - \lambda_{sq,n}} e^{\lambda_{sq,n} \tau} u_{sq,n}(y, k_z, \tau) \langle u_{sq,n}, \mathcal{C}_p \psi_{os,m} \rangle \right)$$



Orr-Sommerfeld:
$$\begin{cases} \mathcal{L} \psi_{os} = \lambda_{os} \psi_{os}, & \psi_{os}(\pm 1) = \psi'_{os}(\pm 1) = 0 \\ \mathcal{S} u_{os} = \lambda_{os} u_{os} - \mathcal{C}_p \psi_{os}, & u_{os}(\pm 1) = 0 \end{cases}$$

Squire:
$$\begin{cases} \lambda_{sq} = - \left(\left(\frac{n\pi}{2} \right)^2 + k_z^2 \right), & \begin{bmatrix} 0 \\ u_{sq} \end{bmatrix} = \begin{bmatrix} 0 \\ \sin \left(\frac{n\pi}{2} (y + 1) \right) \end{bmatrix} \end{cases}$$

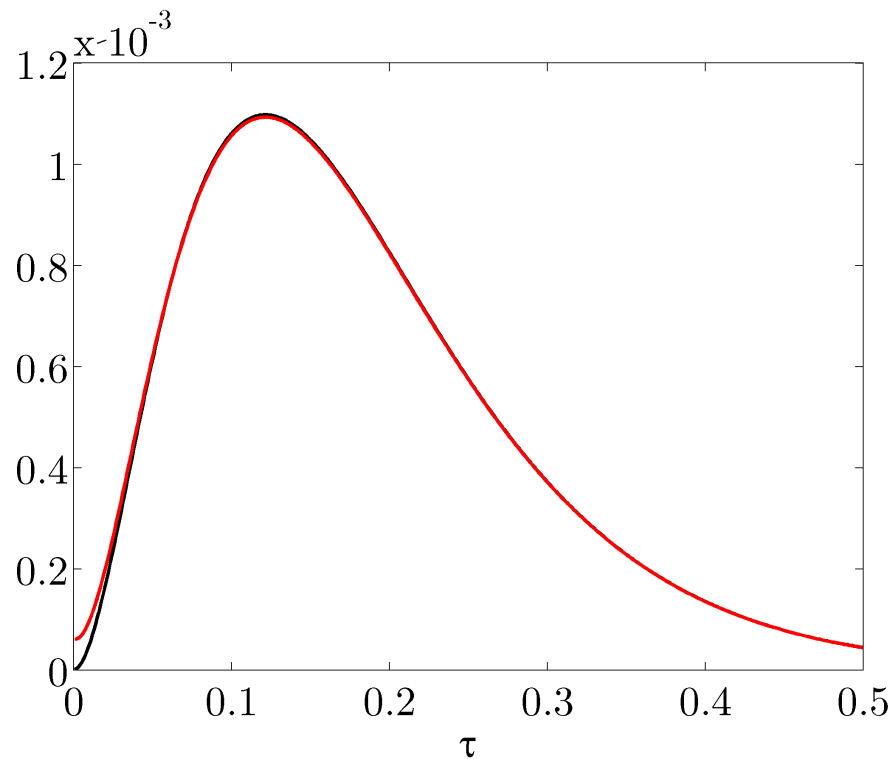
Energy growth

- Worst case energy of u caused by the initial condition in ψ

★ $Re = 1, k_z = 2$

shear-driven flow

(one OS and one Squire mode)



pressure-driven flow

(one OS and two Squire modes)

