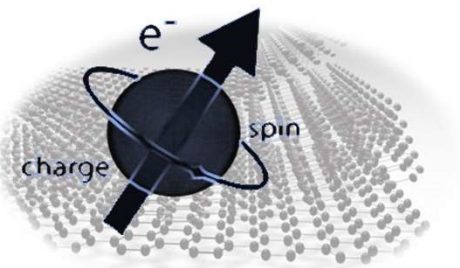


# Spintronics with Topological Insulator Heterostructures

**Saroj P. Dash**

Chalmers University of Technology,  
Sweden

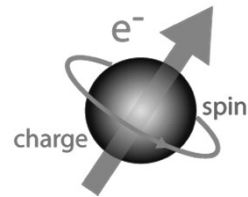


# Research group

MC2

CHALMERS

Quantum Device Physics Laboratory  
Microtechnology and Nanoscience Dept. (MC2)



## Spin transport

- Silicon
- Graphene
- 2D semiconductors
- Topological insulators
- 2D materials heterostructures



Andre



Bogdan



Dmitrii



Anamul



Venkata



Johan



Johannes



Priyamvada



Lennart



Madhushanker



Ravi

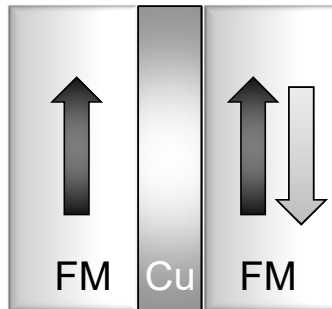


Parham

Christian, Wajid

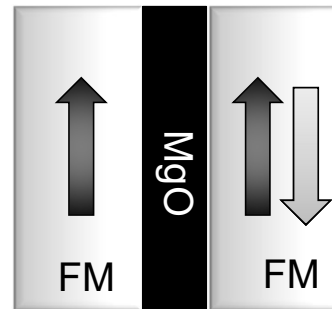
# Spintronic effects

## GMR



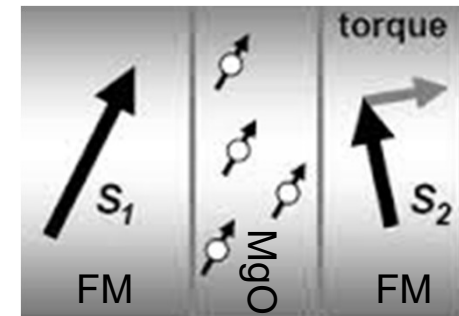
PRL 61, 2472(1988)  
PRB 39, 4828 (1989)

## TMR



PRL 74 3273 (1995)  
Nature Mater 3, 868 (2004)

## STT



Slonczewski, Berger (1996)

P. Grünberg



A. Fert



2007

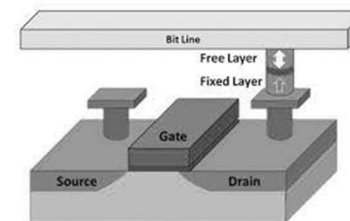
Nobel prize for Physics



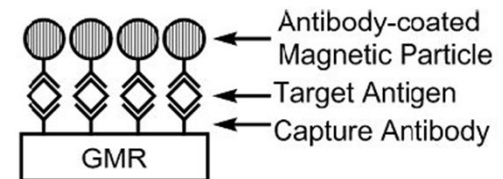
Hard disk



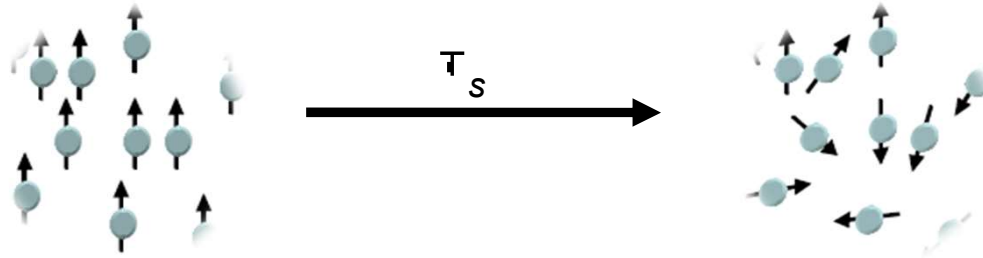
MRAM



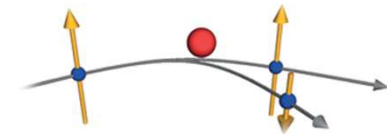
Bio sensors



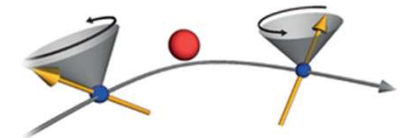
# Spin relaxation



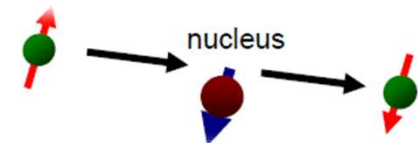
**Elliott-Yafet** - Momentum scattering by phonons or non-magnetic impurities.



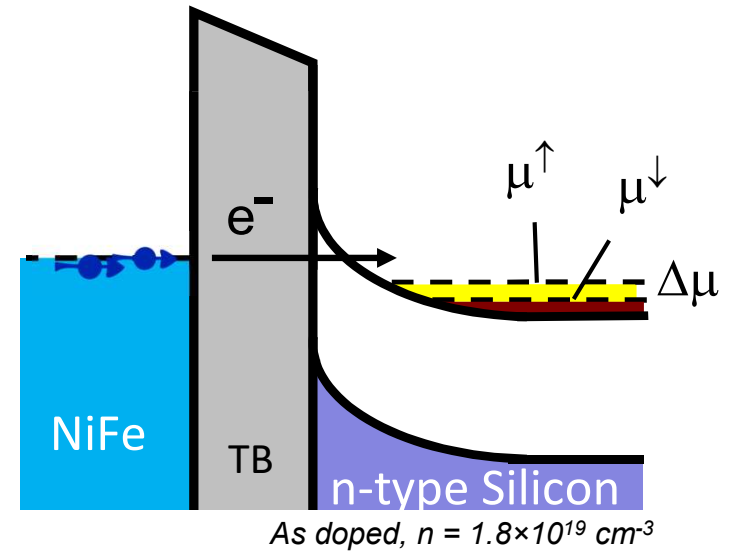
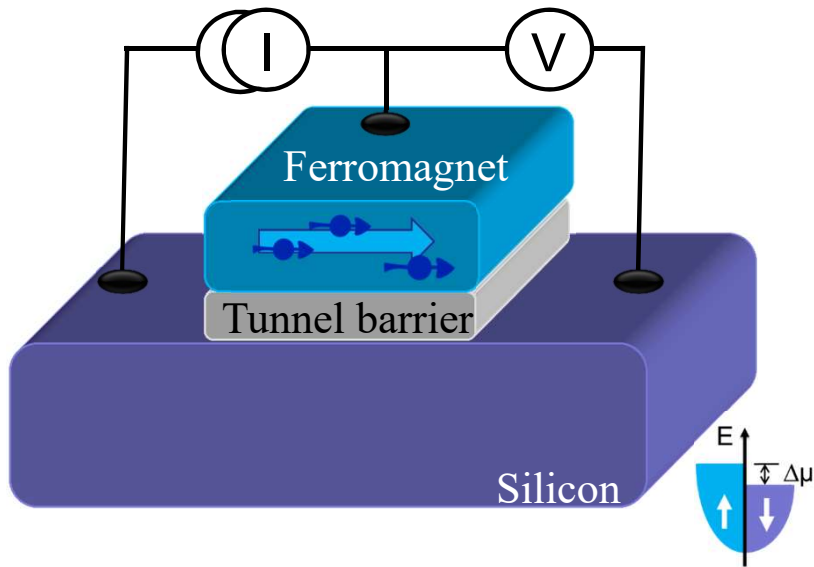
**D'yakonov-Perel'** - Absence of inversion symmetry is equivalent to the presence of an effective magnetic field.



**Hyperfine-interaction** – Interaction between the magnetic momentum of nuclei and electrons.



# Silicon Spintronics

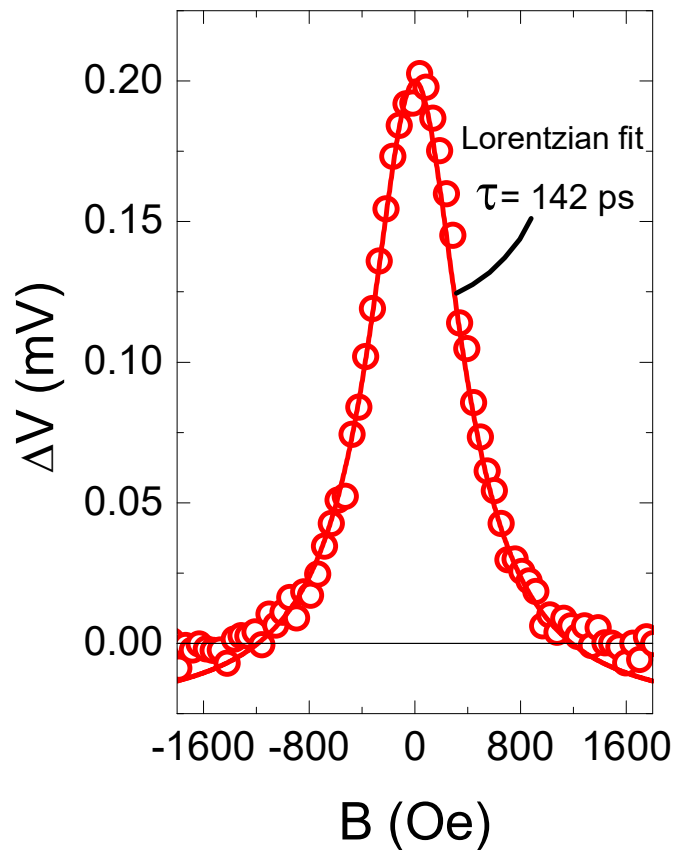


## Longer spin lifetime expected in Silicon

- Low spin-orbit coupling : Inversion symmetry of lattice structure and low atomic number  $Z=14$
- Low hyperfine interaction: Zero nuclear spin for 96 % of Si

S. P. Dash et al. **Nature** 462, 491 (2009)

# Spin injection and detection in Si



S. P. Dash et al. **Nature** 462, 491 (2009)

## Stationary spin-accumulation

$$\Delta\mu(B) = \frac{\Delta\mu(0)}{1 + (\omega_L \cdot \tau_s)^2}$$

## Spin dephasing time (lower limit)

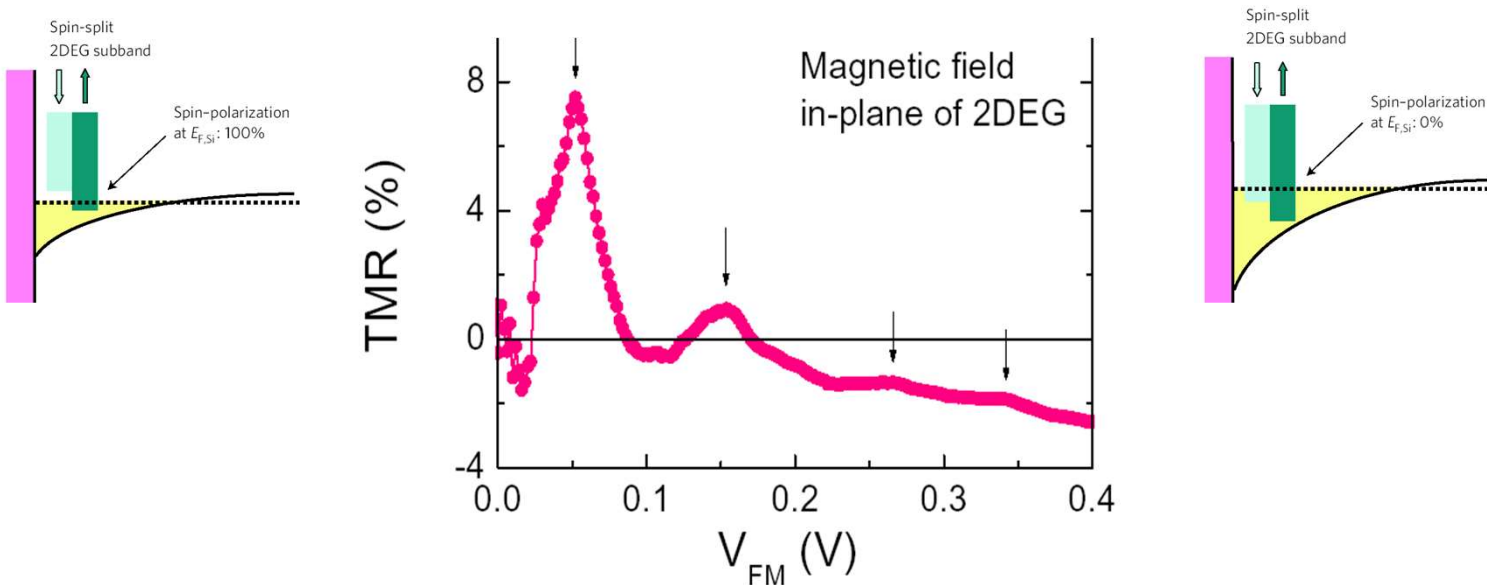
$$\tau_s = \frac{1}{\omega_L} = \frac{\hbar}{g\mu_B(HWHM)} = 140 \text{ ps}$$

## Spin diffusion length (lower limit)

$$L_{SD} = \sqrt{D\tau_s} = 230 \text{ nm}$$

A. Dankert et al. *Sci. Rep.* 3, 3196 (2013)

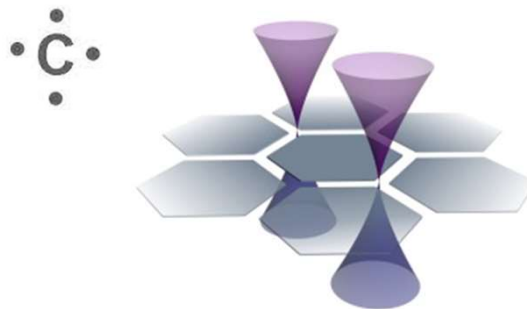
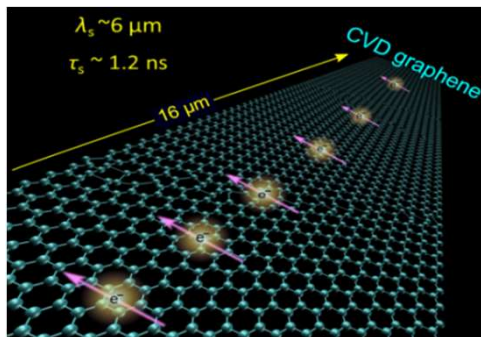
# Spin injection into Si quantum well



- Electrostatic modification of the magnitude of spin polarization in a Si 2DEG
- Detection by means of tunneling to a ferromagnet, producing prominent oscillations of tunnel magnetoresistance of up to 8%.

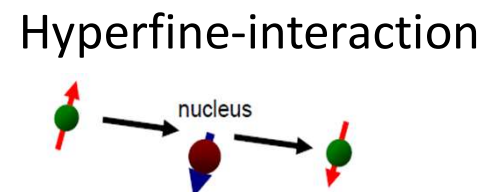
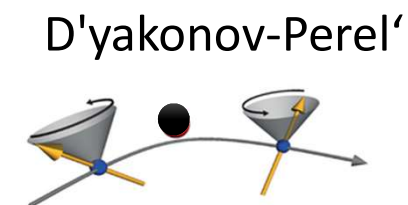
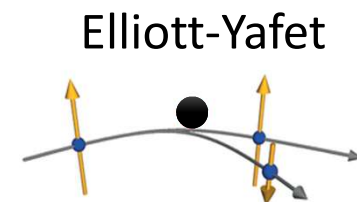
Jansen, Min, Dash, **Nature materials** 9, 133 (2010)

# Graphene Spintronics



## Graphene for long spin coherence

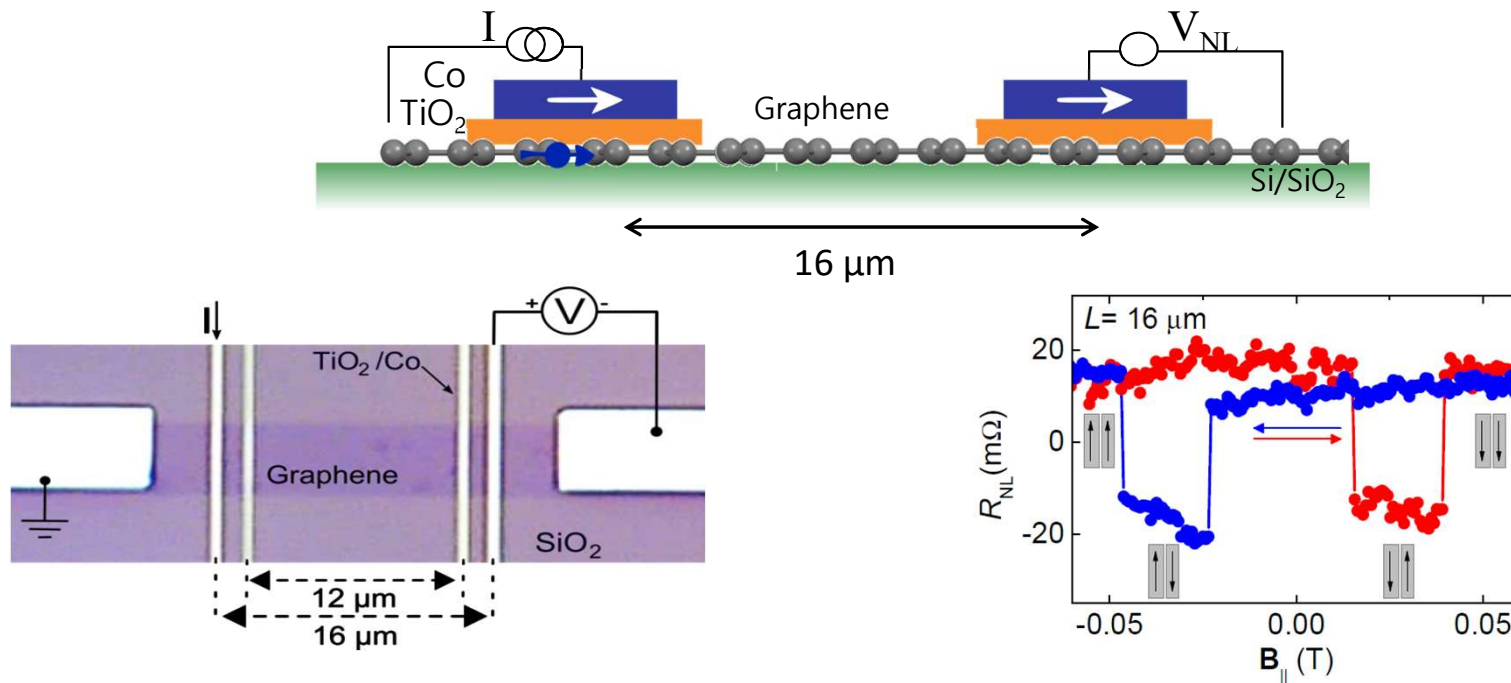
- High mobility
  - Low spin-orbit coupling
  - No hyperfine interaction
- 
- Spin diffusion length 100  $\mu\text{m}$
  - Spin lifetime 1  $\mu\text{s}$



Kamalakar et al., **Nature Communications** 6, 6766 (2015)



# Long distance spin transport in CVD graphene

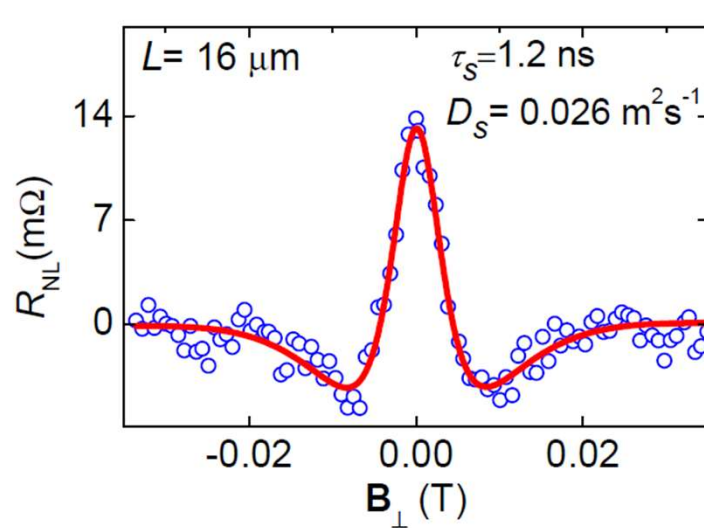
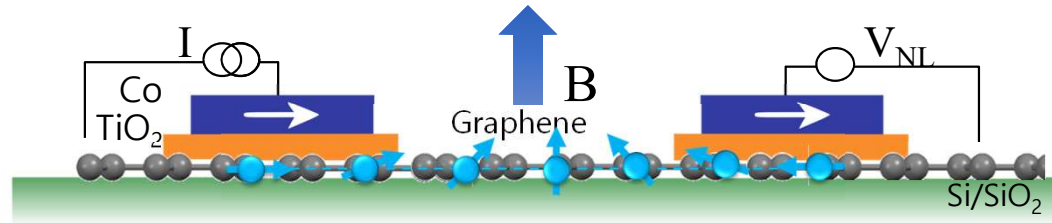


- **Spin transport  $L=16 \mu\text{m}$ ,**
- Higher than previous reports on CVD graphene
- Highest at room temperature - graphene on SiO<sub>2</sub>/Si substrates

Kamalakar et al., **Nature Communications** 6, 6766 (2015)

# Long distance spin transport in CVD graphene

## Spin Precession in Graphene – Hanle effect

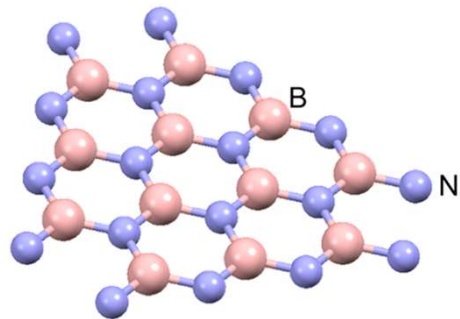
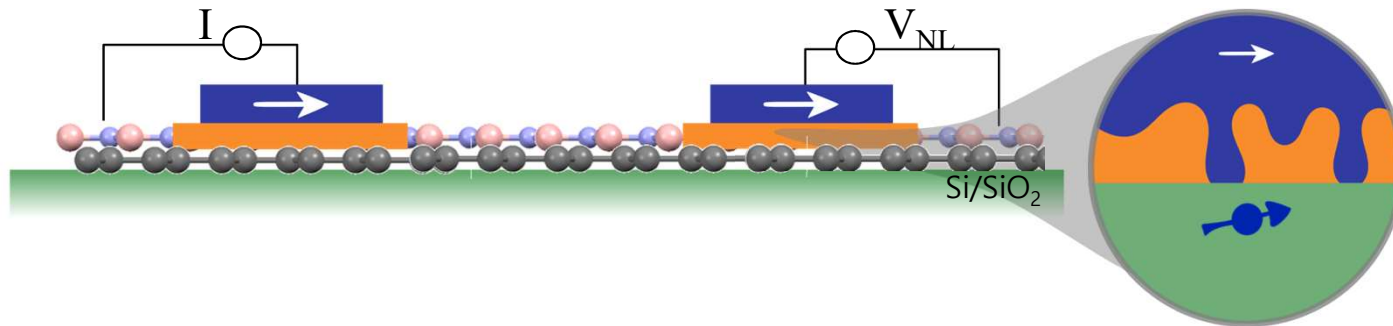


- **Spin parameters  $L=16 \mu\text{m}$ ,  $\tau = 1.2 \text{ ns}$ ,  $\lambda \sim 6 \mu\text{m}$**
- Higher than previous reports on CVD graphene
- Highest at room temperature  
- graphene on  $\text{SiO}_2/\text{Si}$  substrates

Kamalakar et al., **Nature Communications** 6, 6766 (2015)

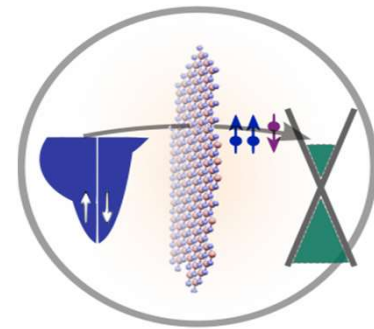
# Graphene/h-BN heterostructures

h-BN tunnel barrier for efficient spin injection into Graphene

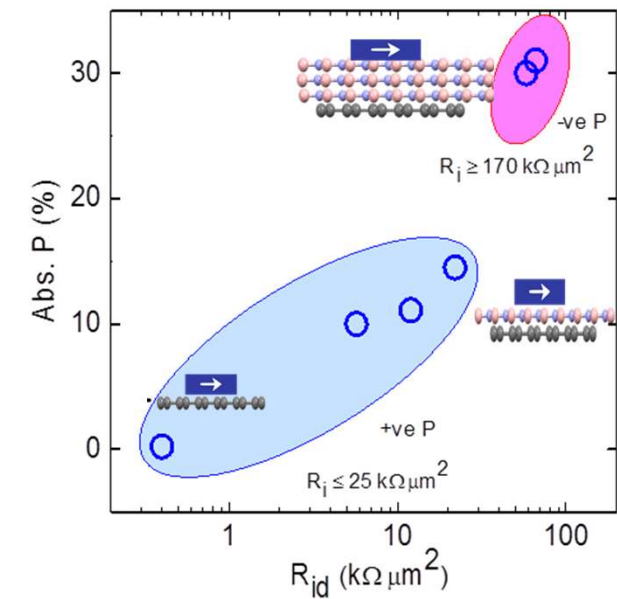
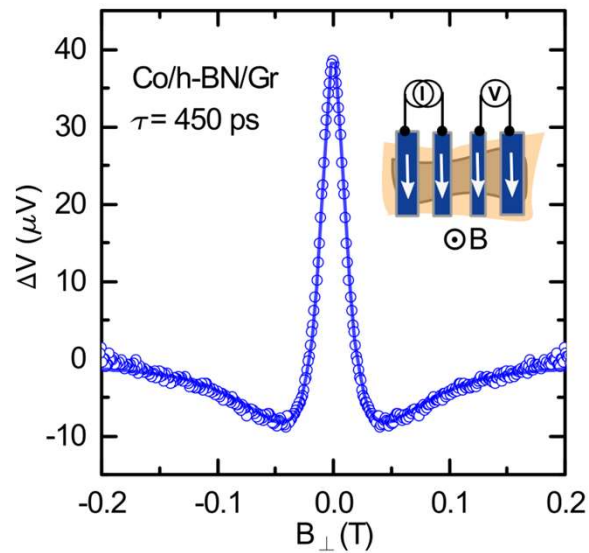
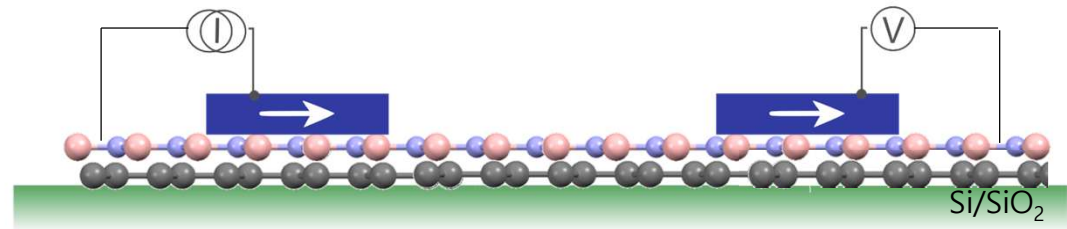
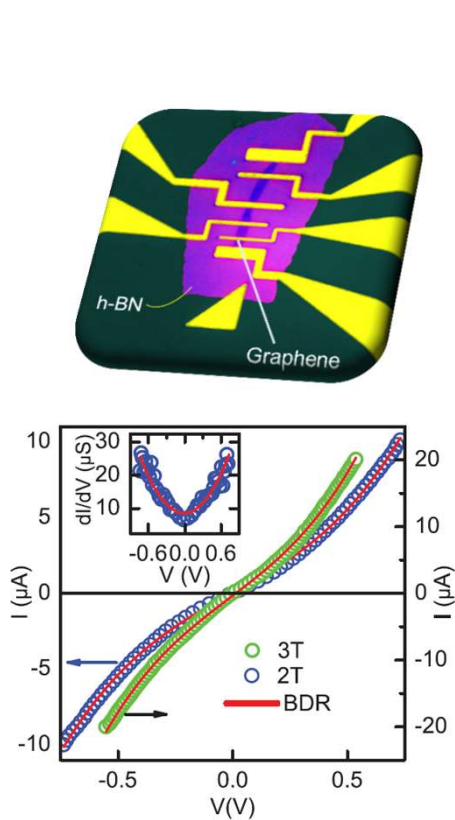


## Hexagonal boron nitride (h-BN)

- Atomically thin insulator  $E_g = 6 \text{ eV}$
- 2D structure  $\rightarrow$  No dangling bonds or interface states.
- Tunnel barrier resistance can be tailored by varying the number of layers



# Graphene/h-BN heterostructures

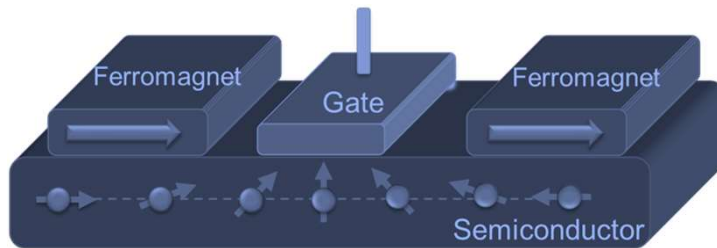


- Improved spin lifetime up to 450 ps
- Larger spin polarization for thicker h-BN barrier

Kamalakar et. al., Scientific Reports 4, 6146 (2014)  
 Scientific Reports 6, 21168 (2016)

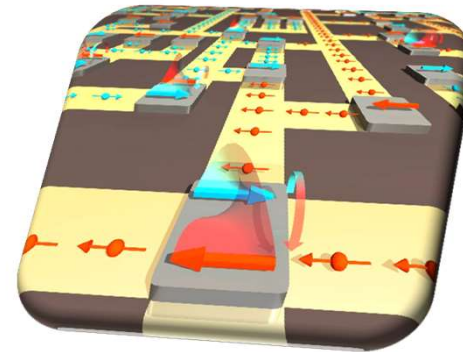
# New perspectives for spintronics

## Spin transistor



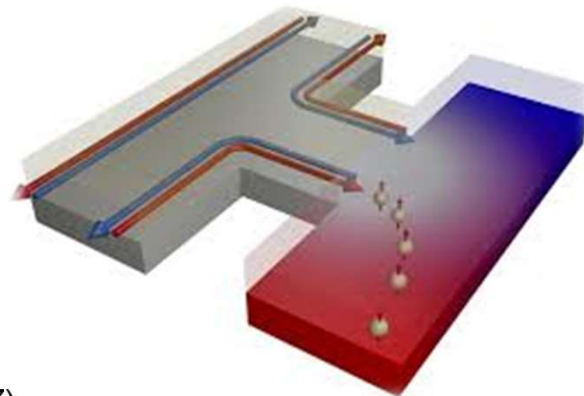
Datta, Das; APL. 56, 665 (1990)

## Spin logic



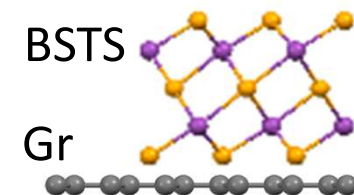
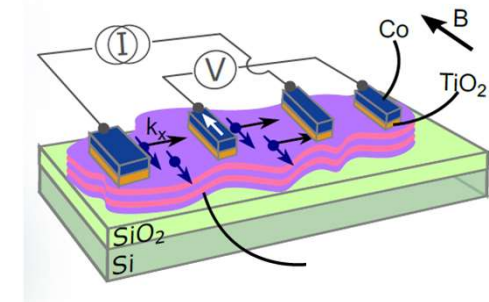
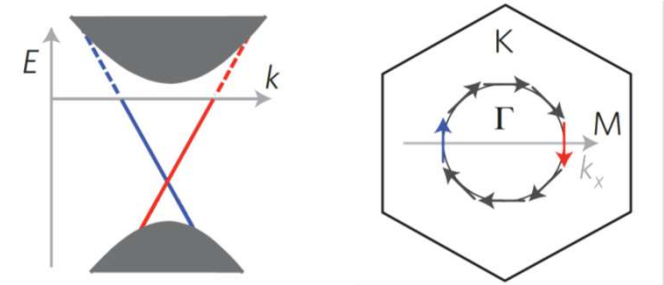
Dery et al. Nature 447, 573 (2007)

## Spin Hall Quantum spin Hall

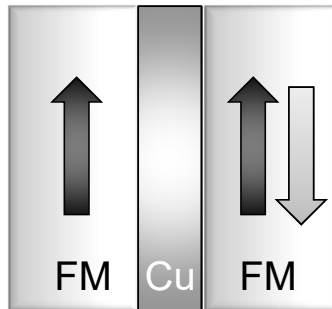


König et al., Science (2007)

- Spin-momentum locking in topological insulators
- Electrical detection of spin-momentum locking up to room temperature
  - $\text{Bi}_2\text{Se}_3$
  - $\text{Bi}_{0.75}\text{Sb}_{1.25}\text{Te}_{0.5}\text{Se}_{2.5}$  (BSTS)
- Dirac material heterostructures of Graphene/Topological insulators

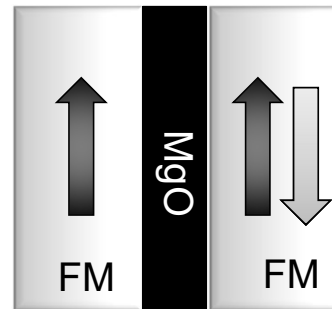


## GMR



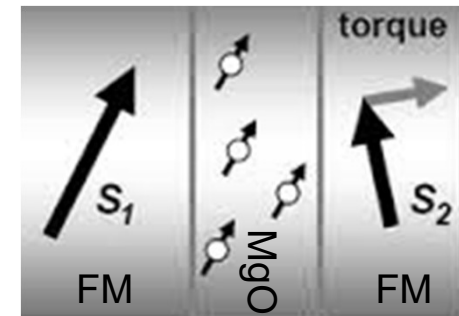
PRL 61, 2472(1988)  
PRB 39, 4828 (1989)

## TMR



PRL 74 3273 (1995)  
Nature Mater 3, 868 (2004)

## STT



Slonczewski, Berger (1996)

P. Grünberg



2007

A. Fert

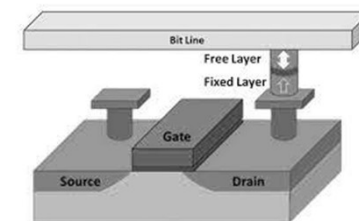


Nobel prize for Physics

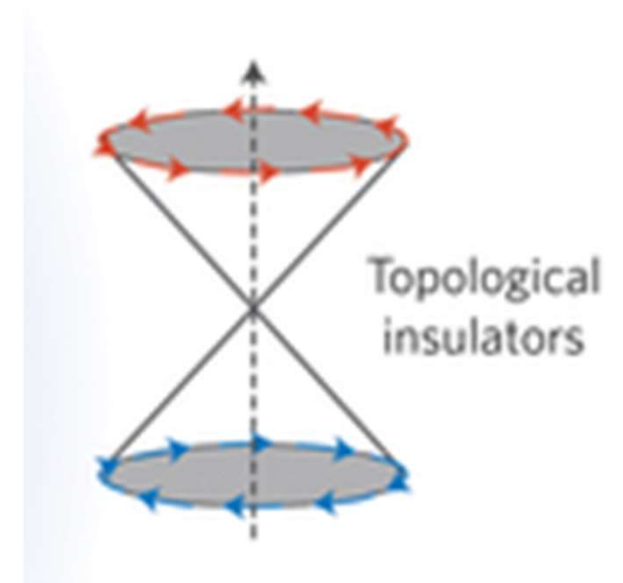
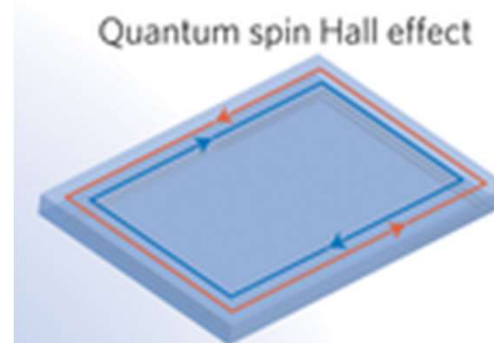
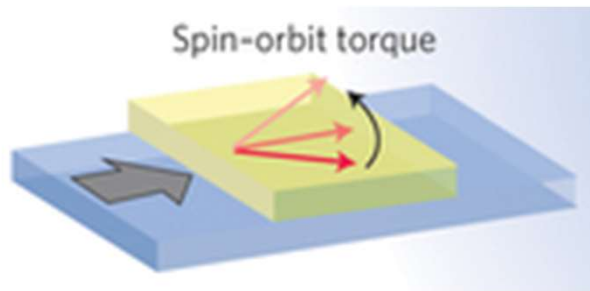
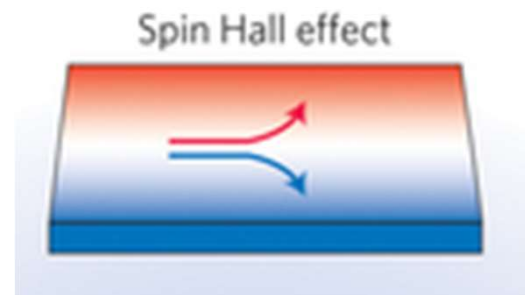
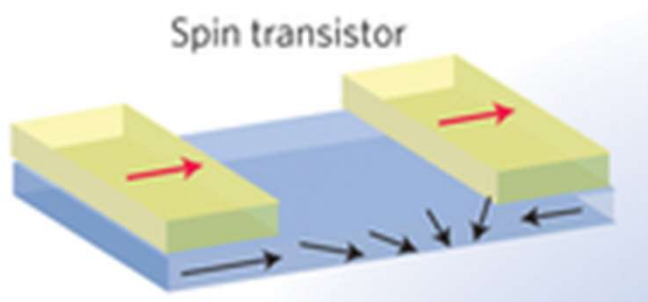
Hard disk



MRAM



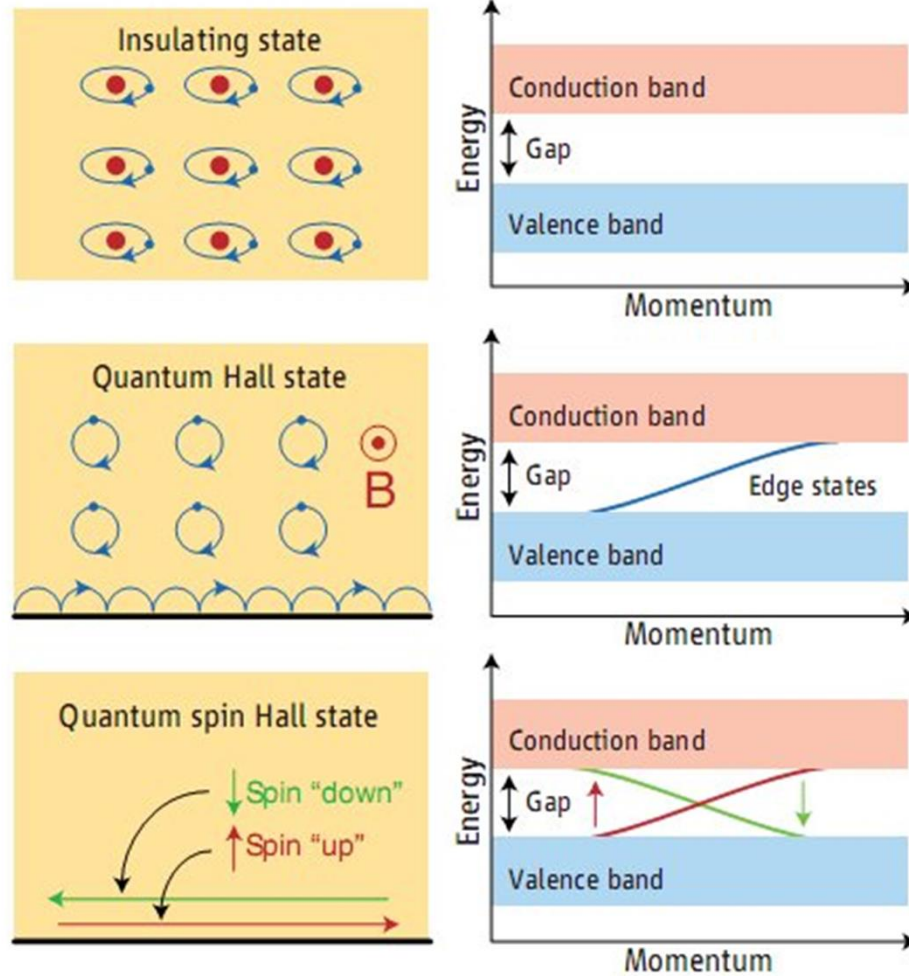
## Spin-orbit coupling



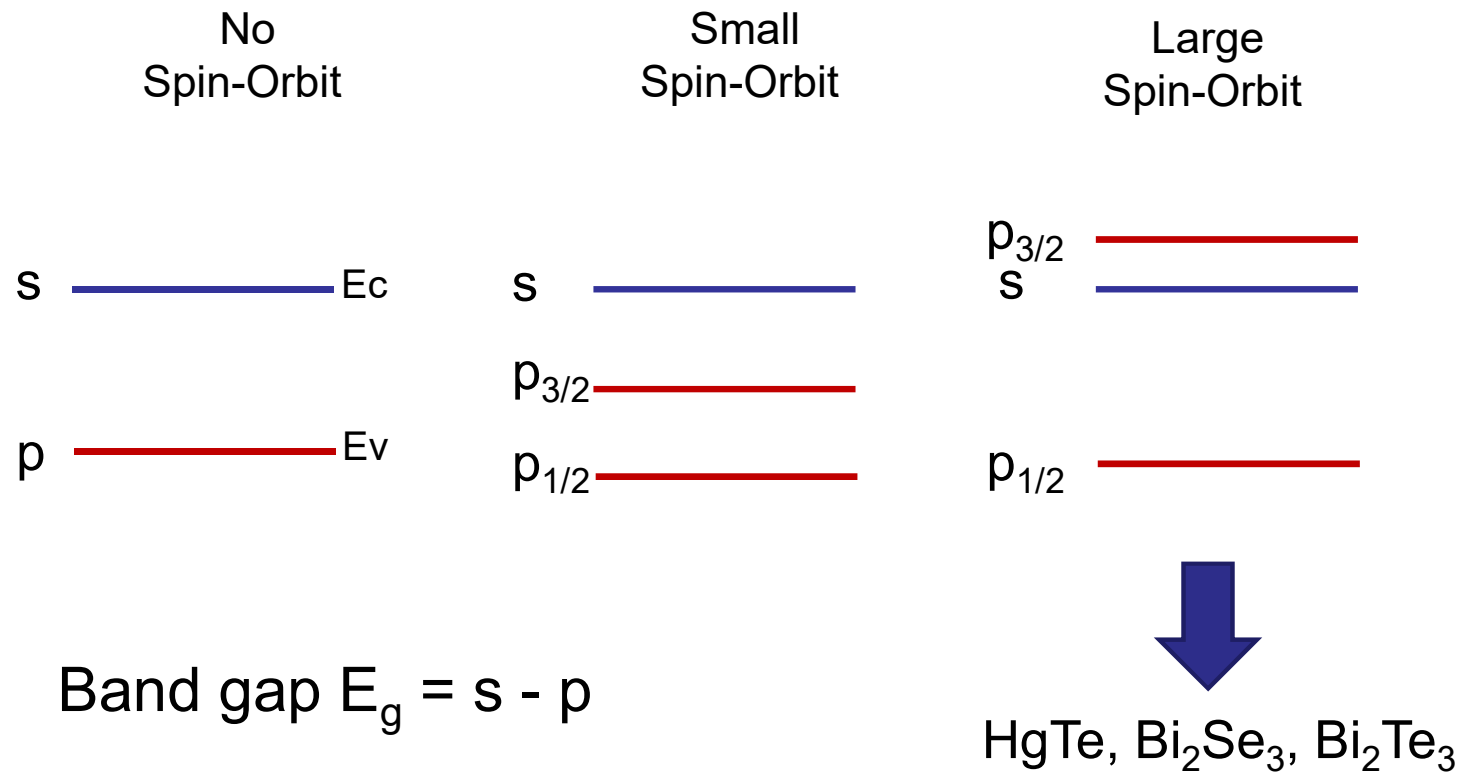


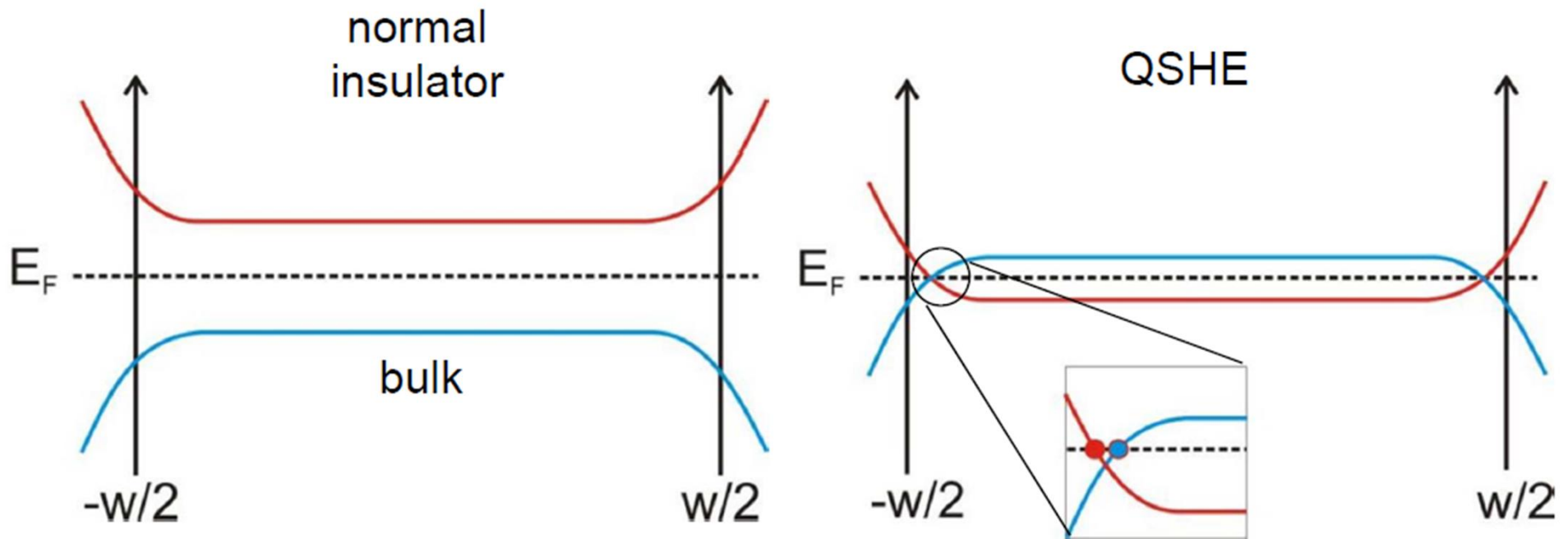
# Topological Insulator

CHALMERS



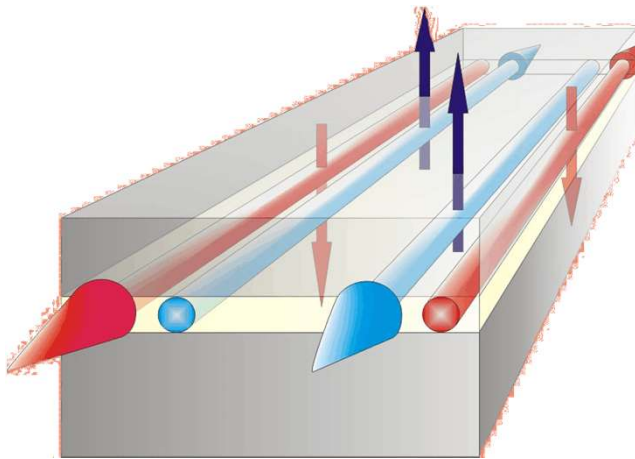
## Spin-Orbit coupling (Band inversion)





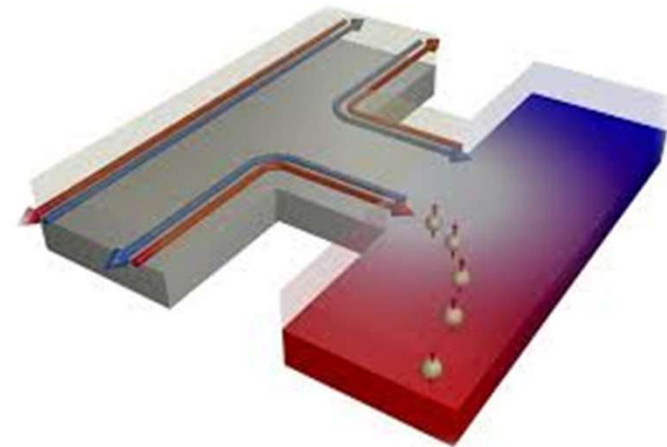
Science 318, 766-770 (2007)

Quantum Spin Hall Insulator State  
in HgTe Quantum Wells

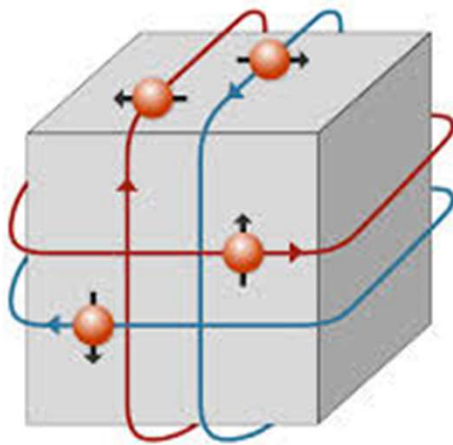
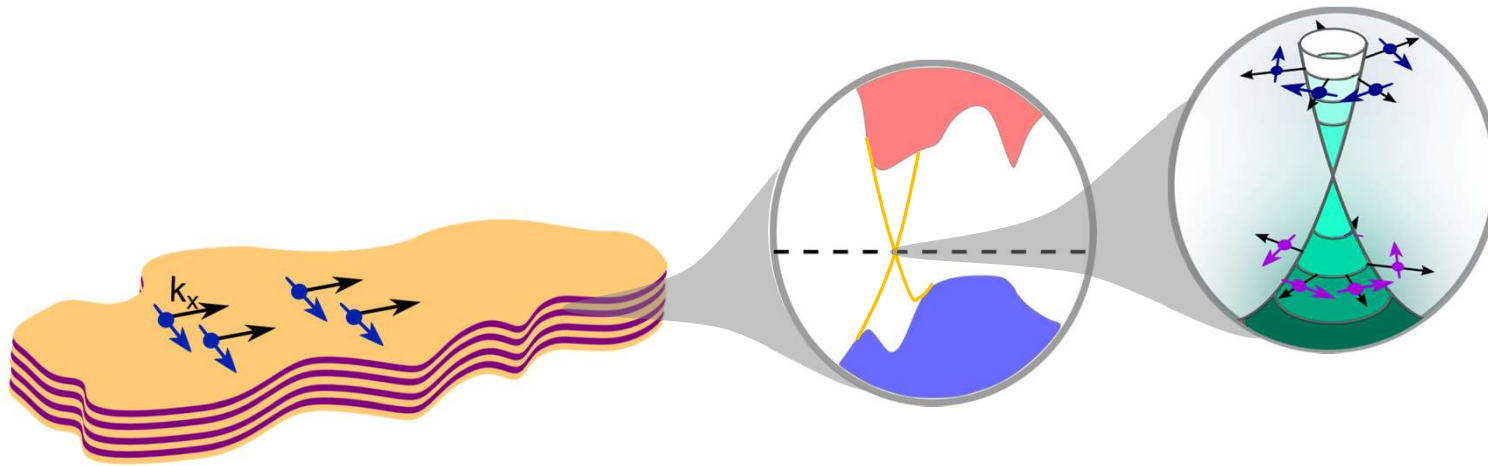


Science 318, 766-770 (2007)

Spin polarization of the  
quantum spin Hall edge states

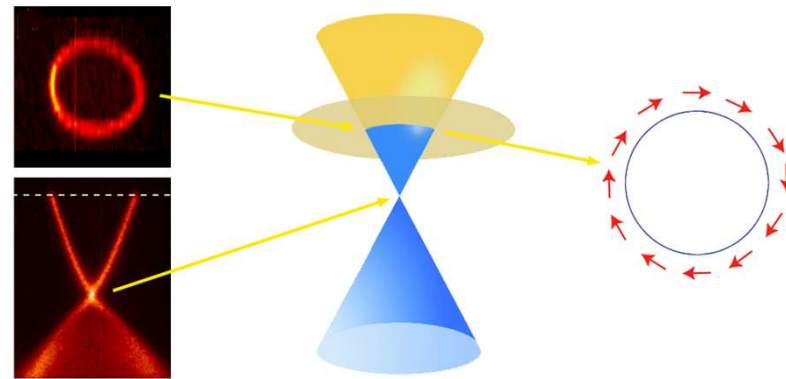
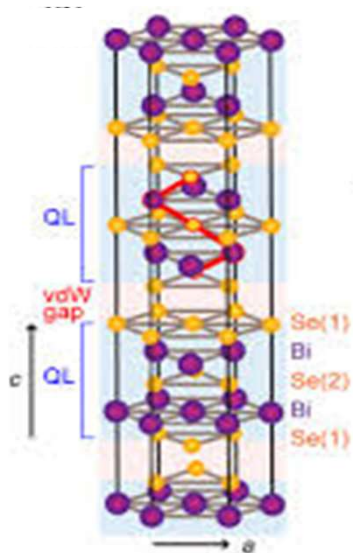


Nature Physics 8, 486-491 (2012)

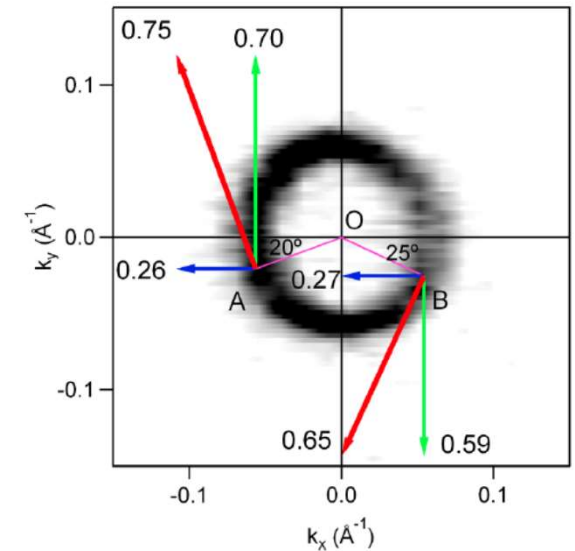


- Insulating bulk band structure
- Dirac state formation at the surface
- Spin and momentum direction are locked (SML)
- Novel measurement techniques possible

## Angle and spin resolved photo emission spectroscopy



Hsieh *et al. Nature* **460**, 1101 (2009)

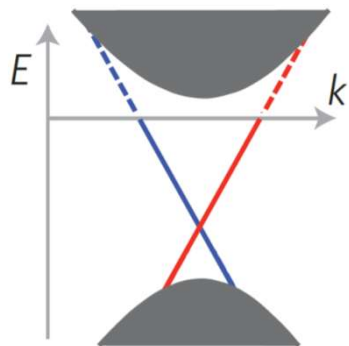
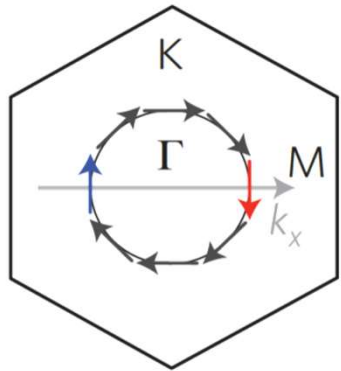


Pan *et al. Phys. Rev. Lett.* **106**, 257004 (2011)

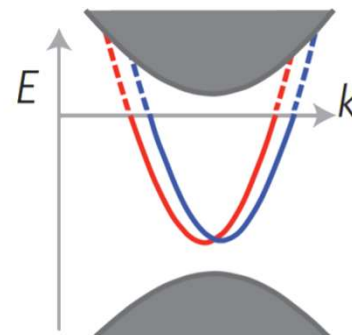
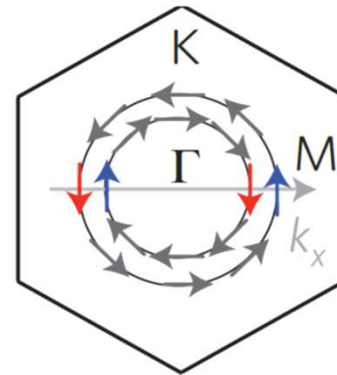
- Bi<sub>2</sub>Se<sub>3</sub>, Bi<sub>2</sub>Te<sub>3</sub>, Sb<sub>2</sub>Te<sub>3</sub> are 3D topological insulators with robust states on surface under room temperature conditions.
- Surface and spin state detected by spin-resolved ARPES

M. Z. Hasan, C. Kane,  
Rev. Mod. Phys. 82, 3045 (2010)

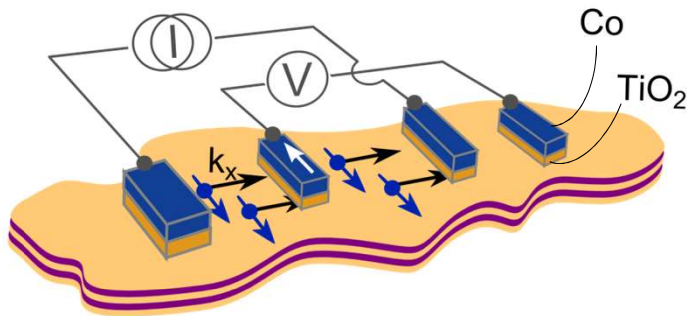
SML



Rashba



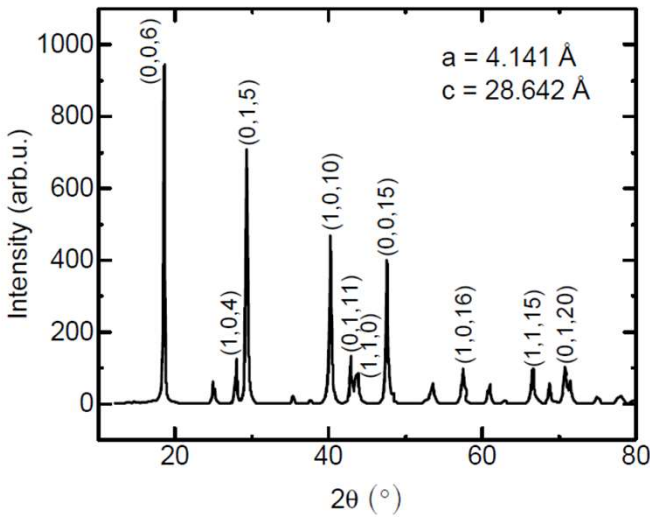
Jozwiak *et al. Nat. Phys.* **9**, 293 (2013)



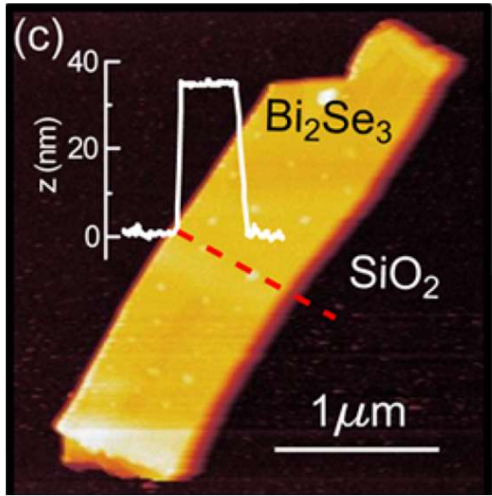
- Applied direct current induces momentum-locked spins
- Detect spins with ferromagnetic tunnel contact
- Net spin polarization can be detected under the contact



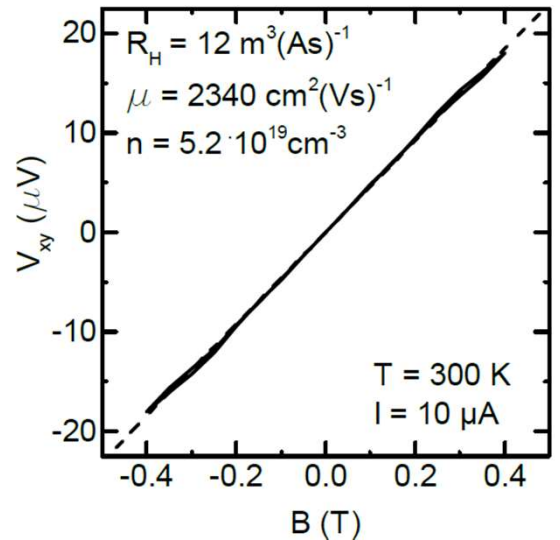
### XRD



### AFM



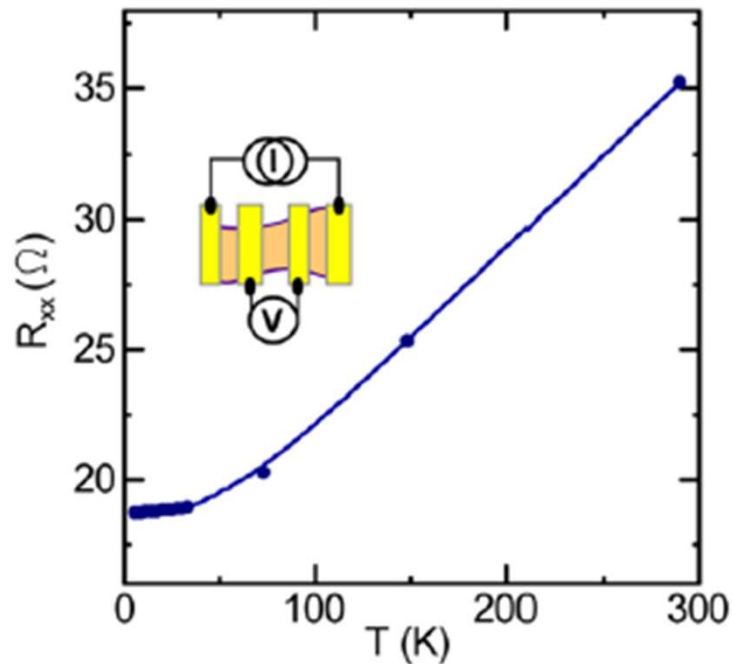
### Hall



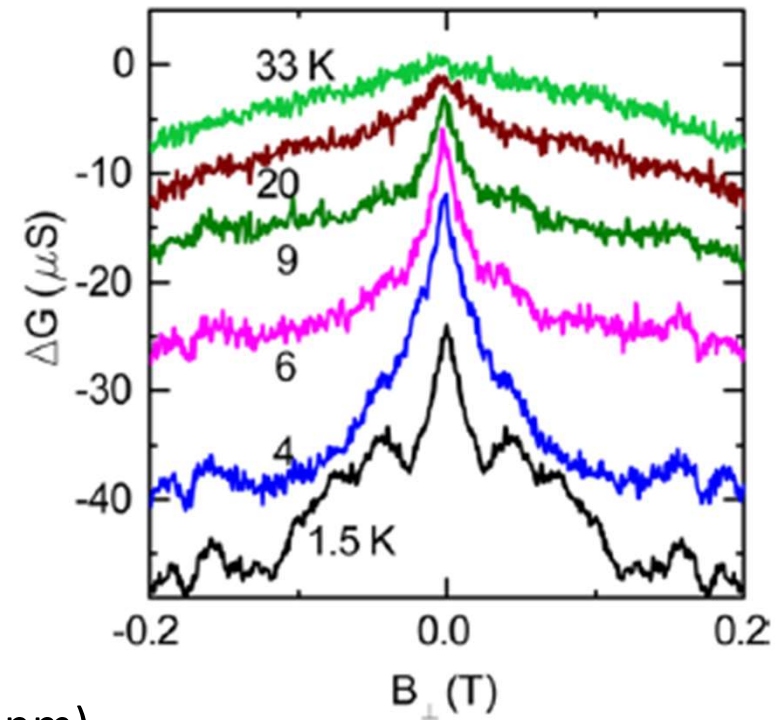
- Crystal grown by Bridgemen method
- $\text{Bi}_2\text{Se}_3$  flakes of 20 – 70 nm thickness
- Mobility  $2000 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$

Dankert, ..., Dash, Nano Letters 15 (12), 7976 (2015)

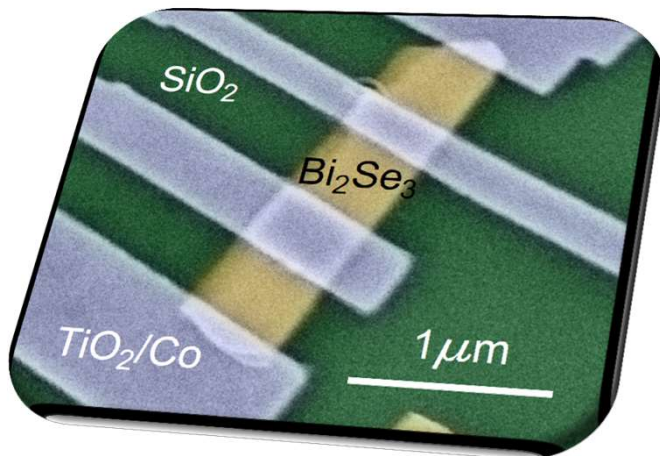
## Resistance



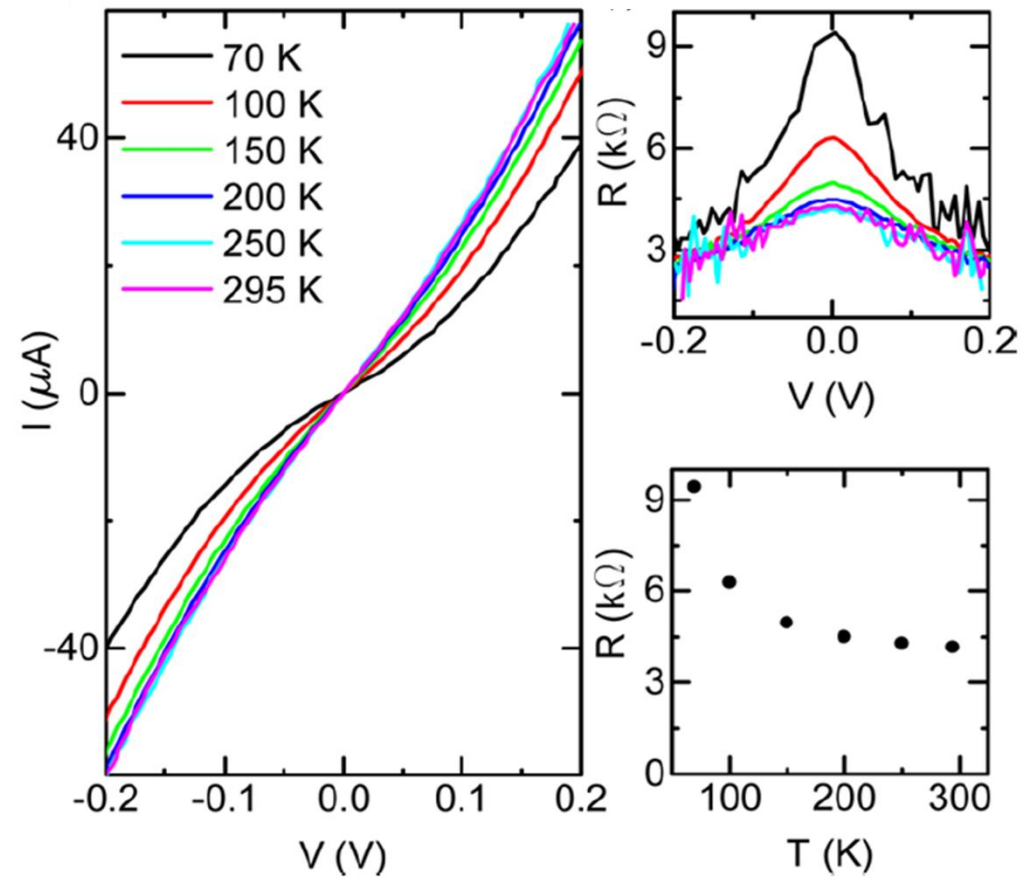
## Weak localization



- Metallic behavior of  $\text{Bi}_2\text{Se}_3$  flakes (20 – 100 nm)
- Weak localization indicates strong spin-orbit coupling in  $\text{Bi}_2\text{Se}_3$

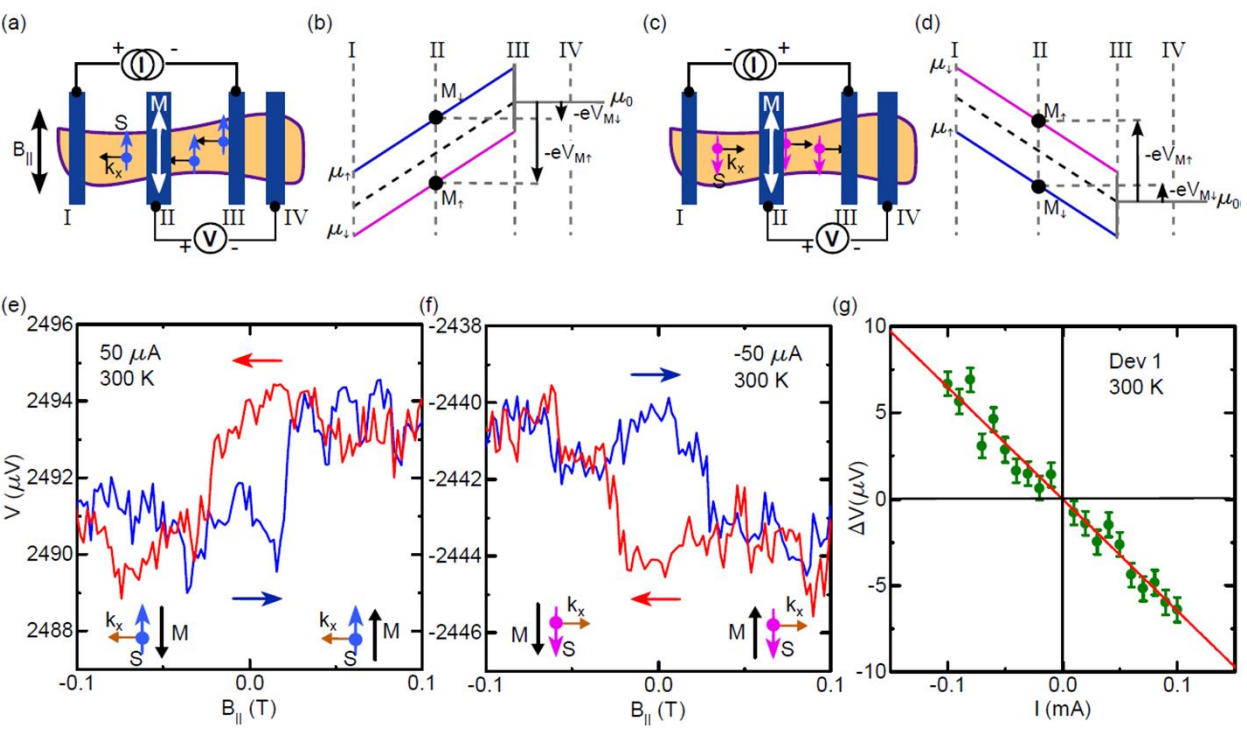


- Exfoliated  $\text{Bi}_2\text{Se}_3$  on  $\text{Si}/\text{SiO}_2$
- $\text{TiO}_2/\text{Co}$  tunnel contacts
- Good tunneling behavior of the contacts

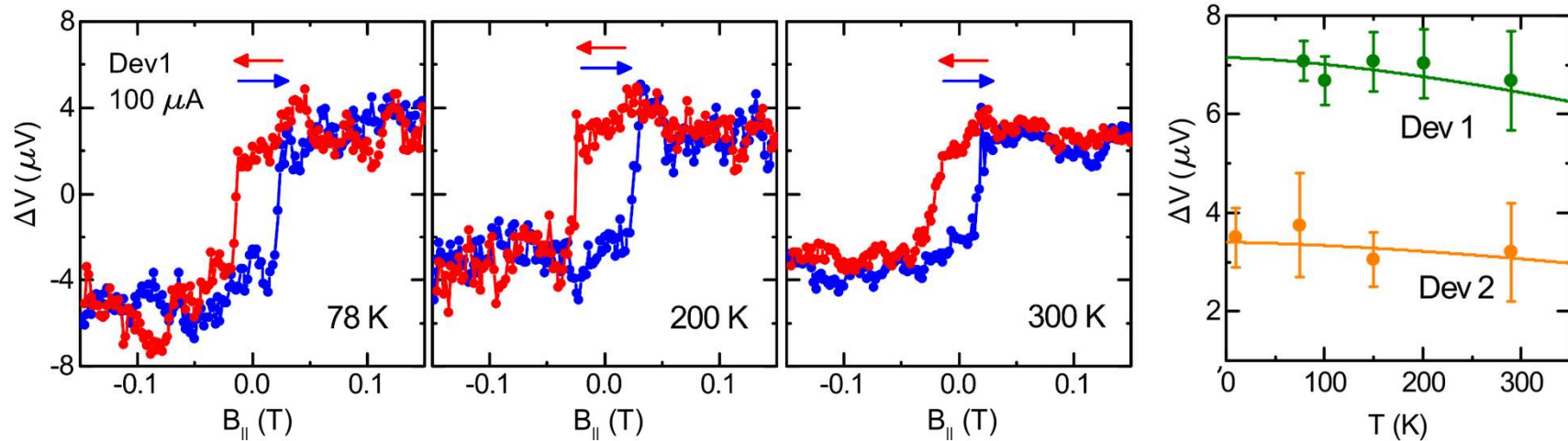


Dankert, ..., Dash, Nano Letters 15 (12), 7976 (2015)

# Detection of spin-momentum locking

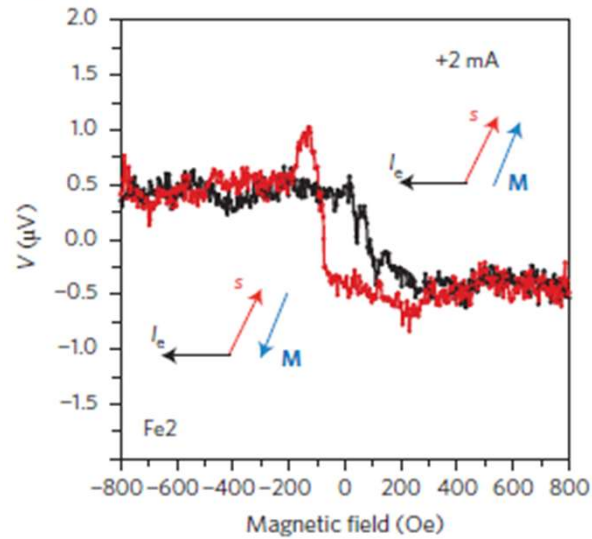


- Potentiometric measurement of spin-polarized surface current
- Spin valve signal between ferromagnet and net surface spins under the contacts
- Spin polarization in surface up to **room temperature**
- Spin-current density  $\propto$  applied current
- Rules out contributions from thermal effects

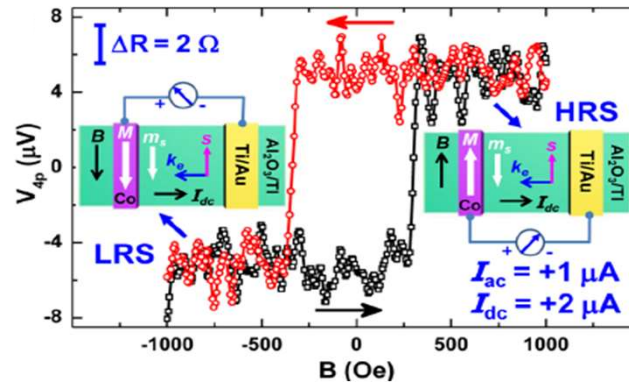


- Very weak temperature dependence  
→ can be influenced by several factors of  $\text{Bi}_2\text{Se}_3$  (bulk/surface contribution, doping, mobility of  $\text{Bi}_2\text{Se}_3$ ) and ferromagnetic tunnel contacts
- $\text{Bi}_2\text{Se}_3$  thickness: Dev1 (40nm) vs. Dev2 (70nm) → Surface origin of signal

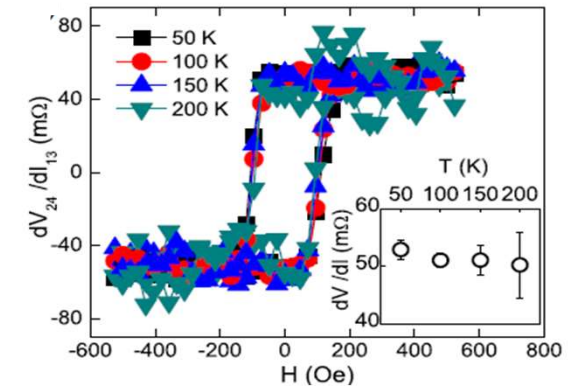
Li et al. Nature Nano 9, 218 (2014)



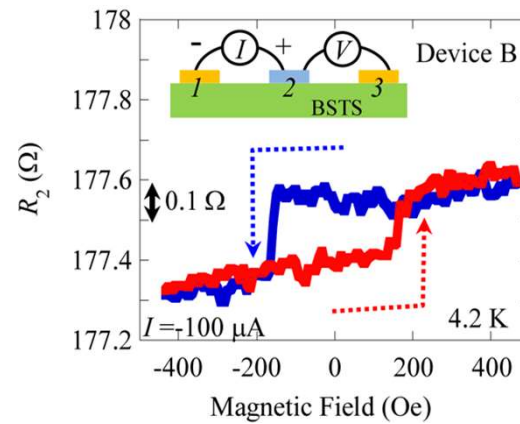
Tang et al. Nano Lett. 14, 5423 (2014)



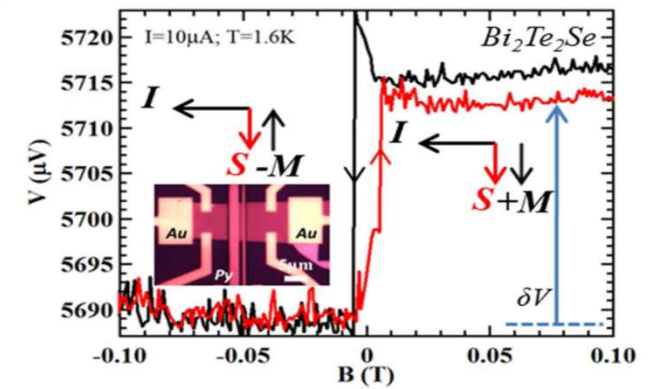
Liu et al. PRB 91, 235437 (2015)



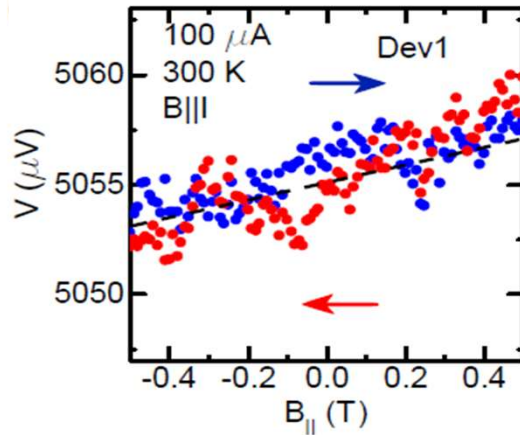
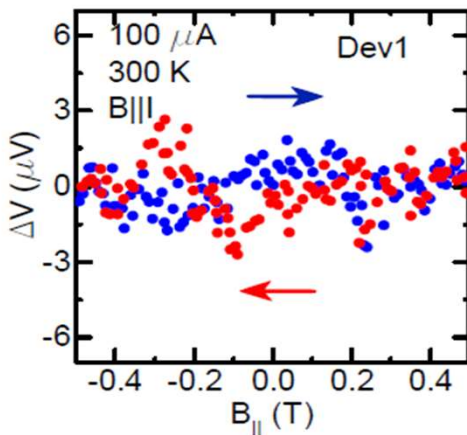
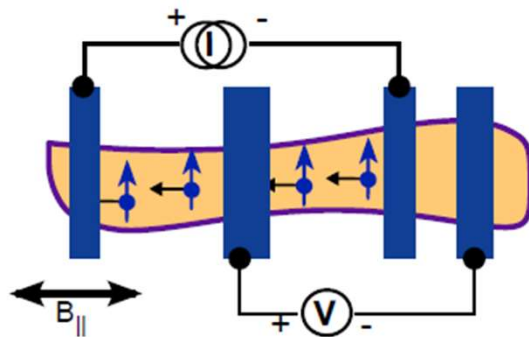
Shiraisi et al. Nano Lett. (2014)



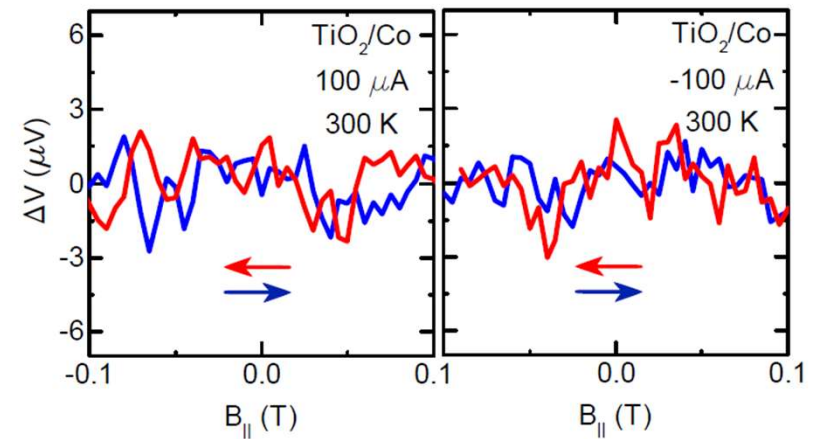
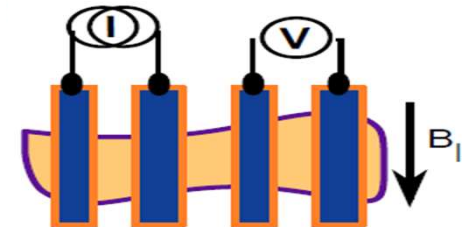
Tian et al. Sci. Rep. 5, 14293 (2015)



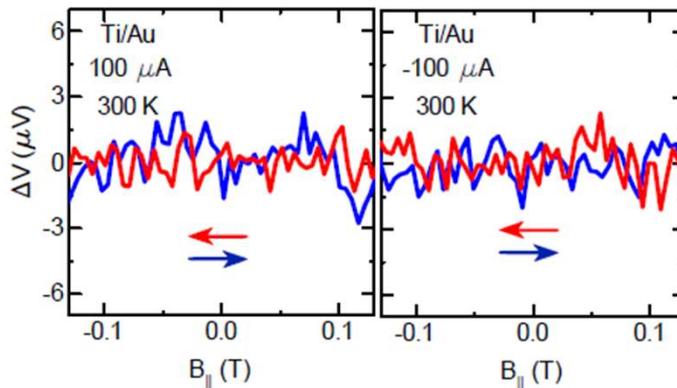
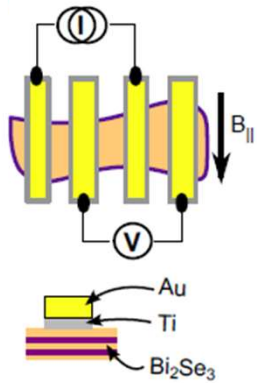
## 90° B<sub>in</sub> Sweep



## Non-local measurement

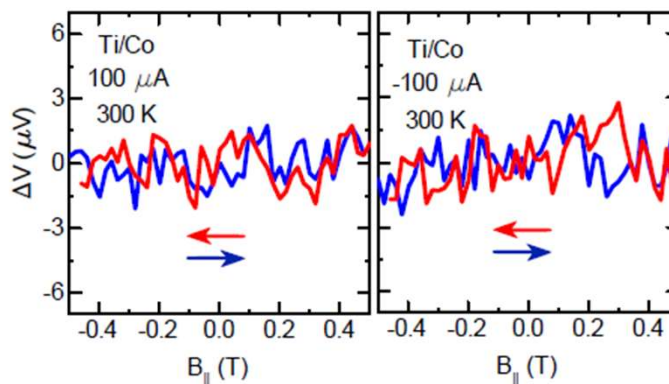
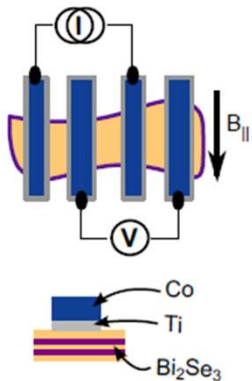


Dankert, ..., Dash, Nano Letters 15 (12), 7976 (2015)



## Non-magnetic Ti/Au contacts

- No signal means, no artifacts from  $\text{Bi}_2\text{Se}_3$



## Magnetic Co contact with non-magnetic Ti interlayer

- No signal means, no artifacts from ferromagnetic electrode



# Summary

CHALMERS

- Observation of spin-momentum locking up to room temperature in  $\text{Bi}_2\text{Se}_3$  and BSTS
- Higher signal in BSTS in comparison to  $\text{Bi}_2\text{Se}_3$
- Stronger temperature dependence of signal in BSTS

