

# Electronic structure, spin transport and magnetic anisotropy of selected cubic Heusler and hexagonal Heusler like alloys

O. Mryasov<sup>1,2,3</sup>

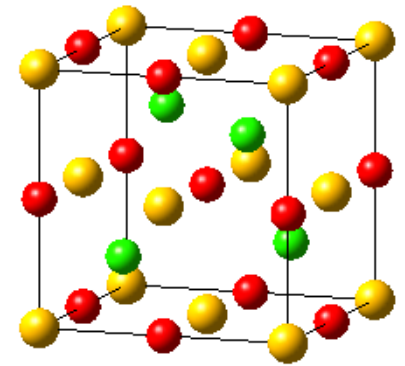
S. Faleev<sup>1,4</sup>, A. Kalitsov<sup>1,3</sup>, J. Barker<sup>1,5</sup>

1. MINT Center, University of Alabama, AL, USA
2. Department of Physics, University of Alabama, USA
3. Western Digital, Advanced Technology, CA, USA
4. IBM, Almaden Research Center, San Jose
5. Tohoku University, IMR, Sendai Japan

# SCOPE : Combination of Properties

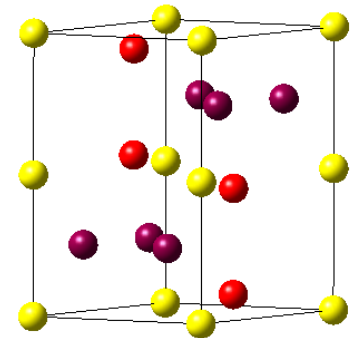
- Structure and composition  $X_2YZ$ ,  $XYZ$ :

- Variety of material
- Mutli-functional properties



- Combination of Properties:

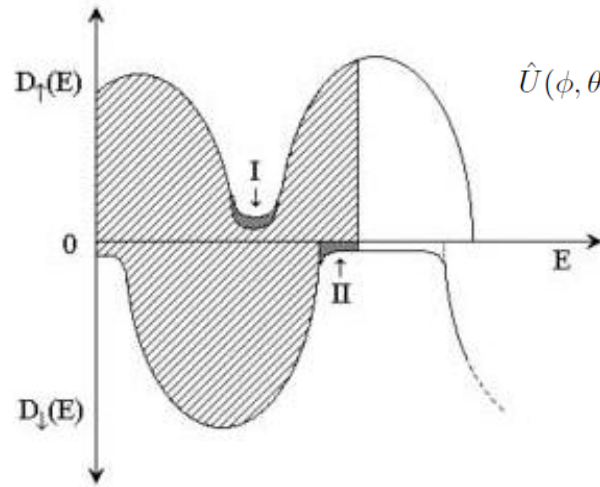
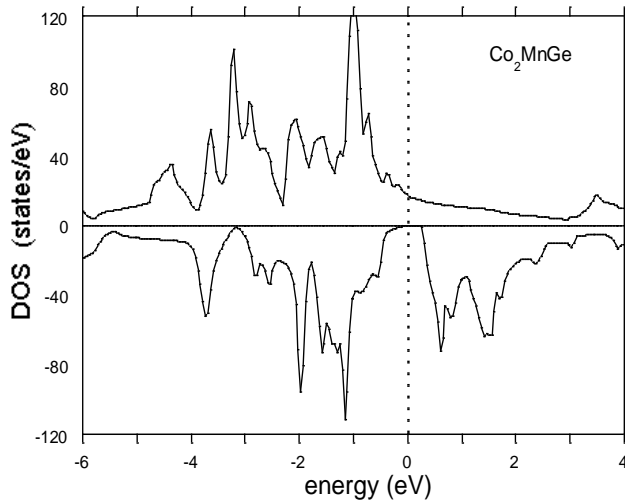
- High spin polarization
- Relatively high Curie point
- Spin dependent transport
- Magnetic anisotropy : bulk or interface



# OUTLINE

- **T<sub>c</sub>** : Comments on Disorder : spin disorder
  - spin mixing - Curie point calculations
  - quaternary alloys strategy
- **E<sub>g</sub>** : Electronic structure: minority band gap
  - fundamental gap theory - alloys design
- **$\rho_{up}/\rho_{dn}$**  : Spin Dependent Transport : GMR
  - Band matching - Q-alloy effects
- **K1** : Hexagonal Heusler like alloys :
  - Magnetic Anisotropy
- Summary and Conclusions

# Minority Band gap and Tc factor



$$\hat{U}(\phi, \theta) = \begin{pmatrix} \cos\left(\frac{\theta}{2}\right)e^{i\theta/2} & \sin\left(\frac{\theta}{2}\right)e^{-i\phi/2} \\ -\sin\left(\frac{\theta}{2}\right)e^{i\phi/2} & \cos\left(\frac{\theta}{2}\right)e^{-i\phi/2} \end{pmatrix}$$

Prof. Hono

- Increase Minority Gap

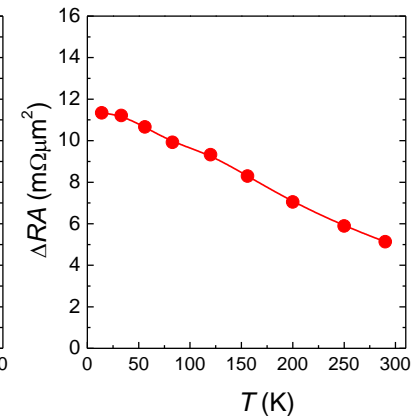
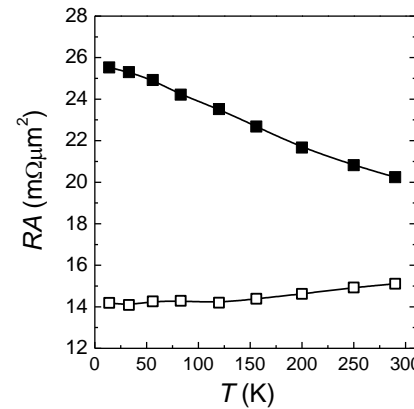
CFGG = Co<sub>2</sub>Fe(Ge<sub>0.5</sub>Ga<sub>0.5</sub>)

CMGG = Co<sub>2</sub>Mn(Ge<sub>0.75</sub>Ga<sub>0.25</sub>)

CFAS = Co<sub>2</sub>Fe(Al<sub>0.5</sub>Si<sub>0.5</sub>)

CMFS = Co<sub>2</sub>(Mn-Fe)Si

CMFG = Co<sub>2</sub>(Mn-Fe)Ge



- Increase Tc

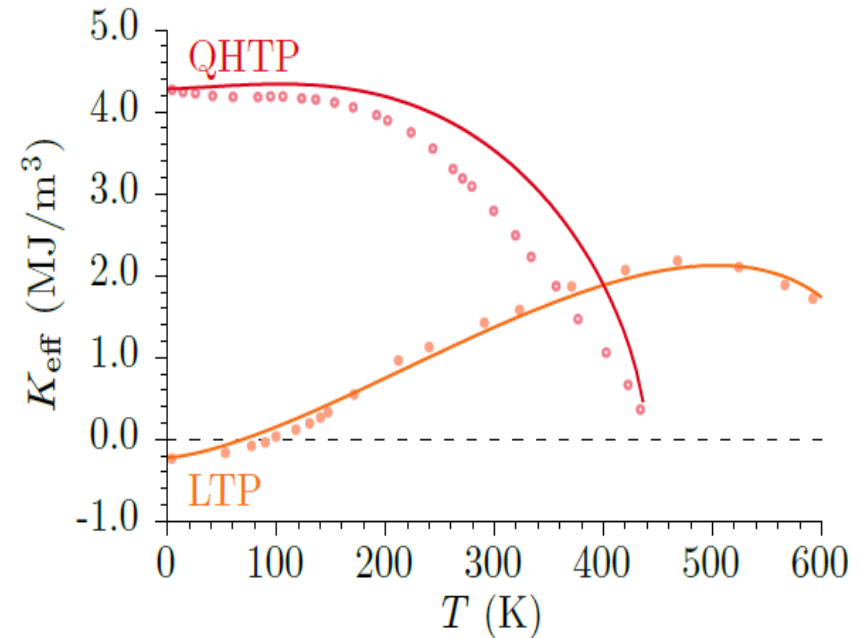
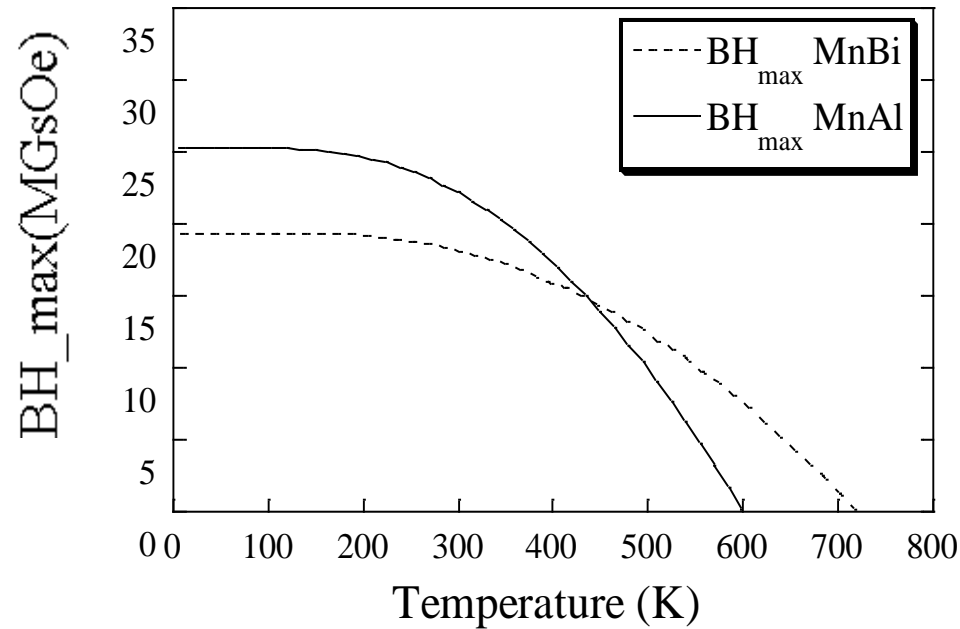
- Co<sub>2</sub>FeGe

- Co<sub>2</sub>FeGa larger Tc

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# How to calculate $M(T)$ , $K(T)$ , $P(T)$ ?



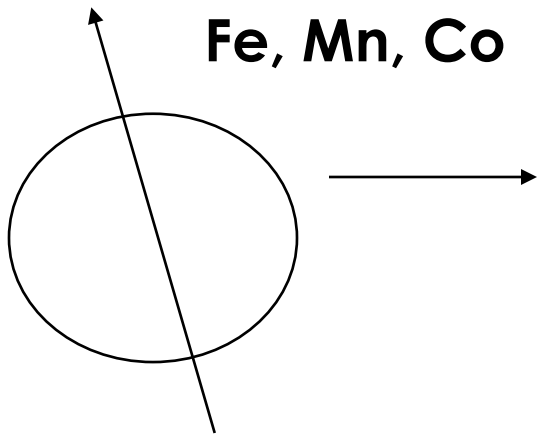
- Effective Spin Hamiltonian approach

O. N. Mryasv *et.al.*, *EuroPhysics Letters*, **69**(5), p.805 (2005)

- statistical simulations/theory
- material specific parameterization

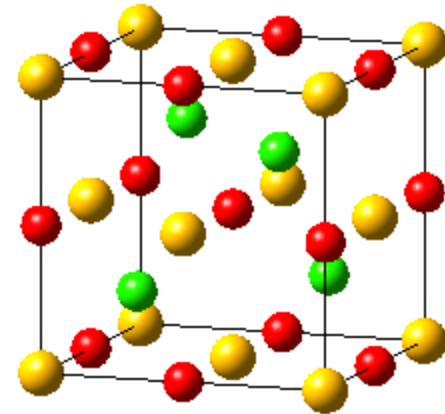
# Spin Hamiltonian: Microscopic Definition

- Two terms in the effective potential variation:



$$\Delta V_i^{ex} = B_i \delta \vec{e}_i \vec{\sigma}$$

$$\Delta V_i^{so} = \frac{1}{2} \sum_l \xi_l^i \vec{L} \delta \vec{e} \vec{\sigma}$$



- Contributions to the total energy:  $\delta E = \delta E^{EX} + \delta E^{DM} + \delta E^{MAE}$

$$E^{EX} = - \sum J_{ij} \vec{e}_i \vec{e}_j \quad E_{DM} = \sum D_{ij} [\vec{e}_i \times \vec{e}_j] \quad E^{MAE} = \sum \varepsilon_{ij}$$

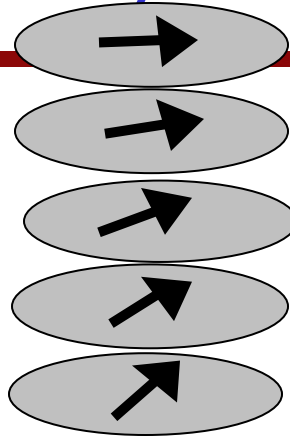
$$E^{EX} = - \sum J_{ij} \mathbf{e}_i \mathbf{e}_j$$

$$E^{MAE} = - \sum k_{\ddot{u}}^{(0)} \mathbf{e}_i^z \mathbf{e}_i^z - \sum k_{ij}^{(2)} \mathbf{e}_i^z \mathbf{e}_j^z$$

# Generalized constrained density functional theory

$$E_{ex} = AVq^2 \quad \bullet \text{ Spin spiral excitations}$$

$$A(x, T) = A(0, T) [1 - \lim_{q \rightarrow 0} \langle\langle \Delta E_{ss}(x, q) \rangle_x \rangle_T / \langle E_{ss}(0, q) \rangle_T]$$



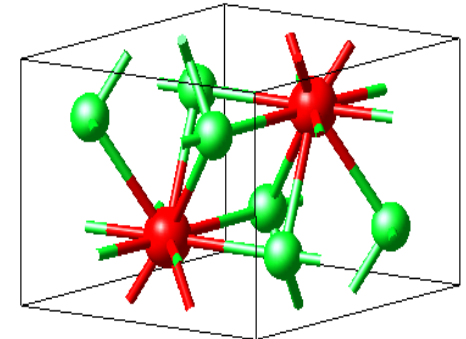
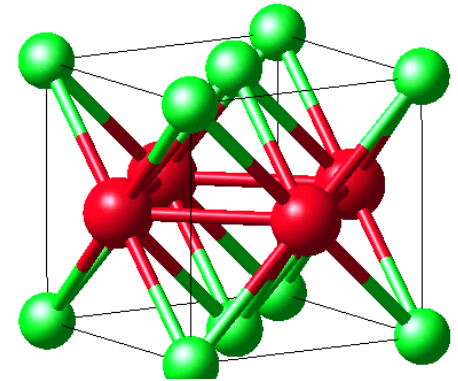
O. Mryasov et.al Phys. Rev. B. 45, 12330 (1992)

- Constrained DFT calculations

$$\left[ -\frac{\hbar^2}{2m} \nabla^2 + V_{eff}(\vec{r}) \right] \varphi_i(\vec{r}) = \varepsilon_i \varphi_i(\vec{r})$$

$$\vec{B} = \left( \frac{\partial E_{xc}[\rho, m]}{\partial m} \frac{\vec{m}}{m} \right)$$

$$H_{so} = -\xi LS$$



$$E_{CLDA}[\rho, B, h_{\perp}] = E_{LDA}[\rho, B] + E_{Const}[\rho, h_{\perp}]$$



# Multi-sub-lattice mean field

	Tc (Experiment)	Tc (Theory)
	(K)	(K)
Fe (bcc)	1040	1080
Co (fcc)	1400	1533
Co <sub>2</sub> FeGe	1000	1062
Co <sub>2</sub> MnGe	905	867
Co <sub>2</sub> FeSi	1120	1047
Co <sub>2</sub> MnSi	985	963
Co <sub>2</sub> FeAl	1000	1298
Co <sub>2</sub> MnAl	693	590

- Y = Fe vs. Mn - approach to increase Tc
- Q: How this changes band gaps

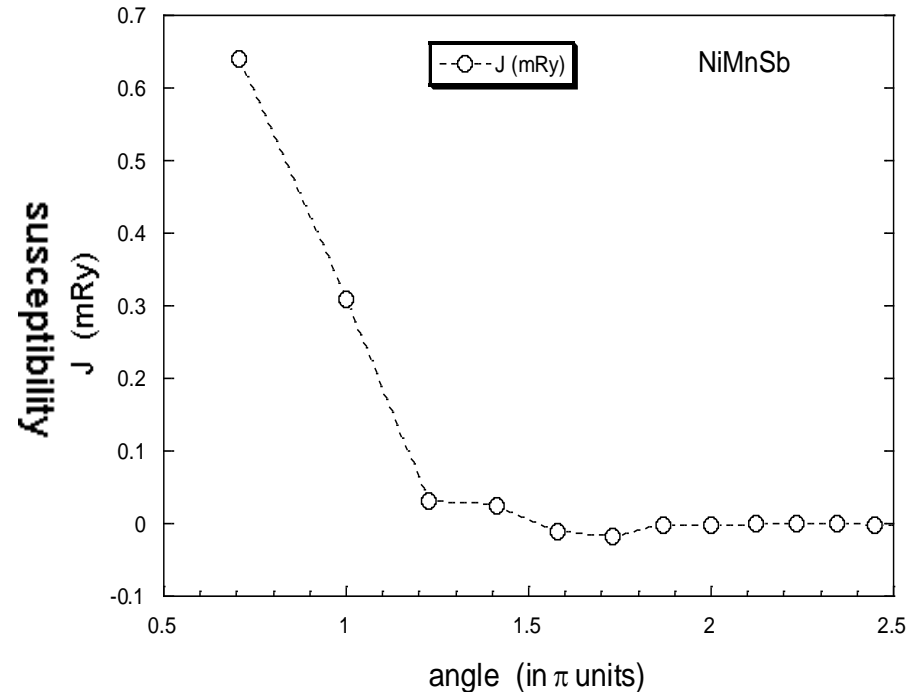
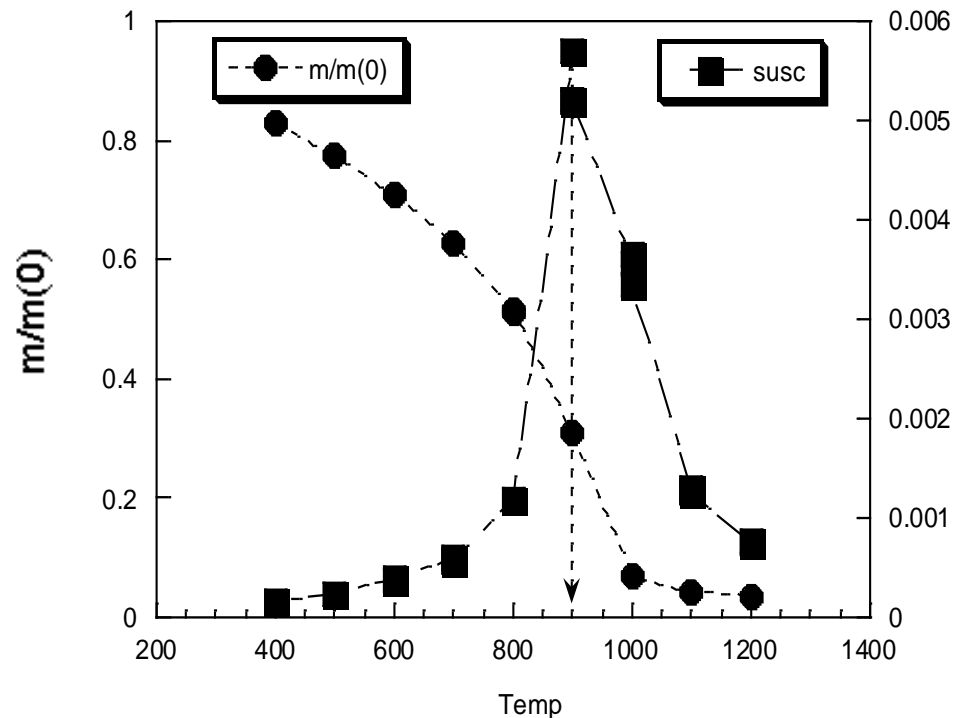
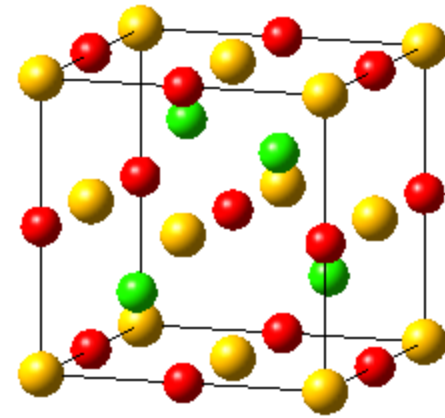
# Tc calculations beyond mean field

Mean field theory for Two sub-lattice magnets:

$$T_C = \frac{1}{2}(T_{TT} + T_{RR}) + \sqrt{\frac{1}{4}(T_{TT} - T_{RR})^2 + T_{RT}^2}$$

NiMnSb alloy test

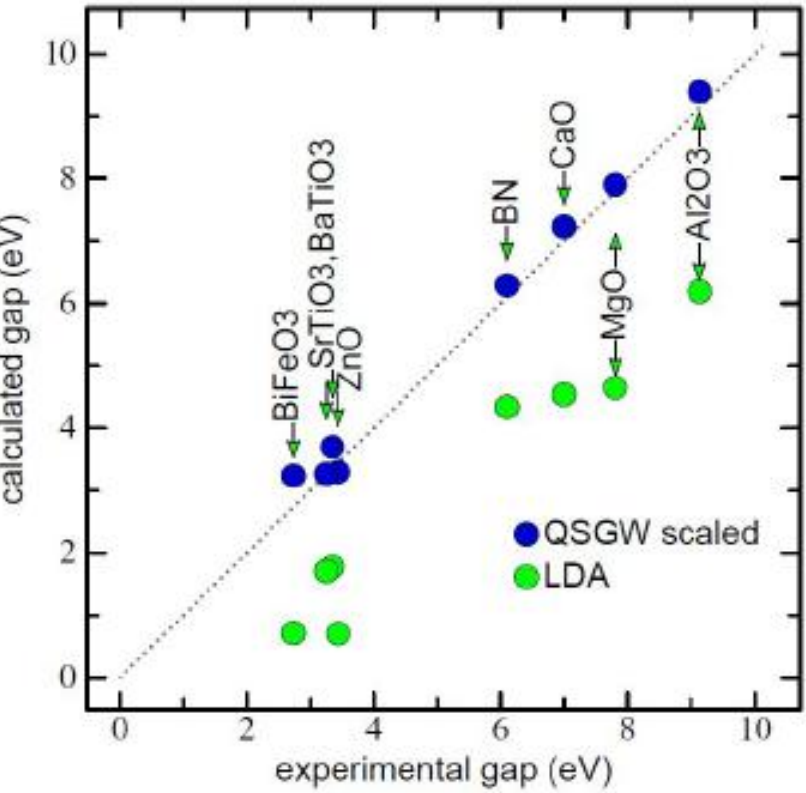
Experimental Tc about 780 K



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# Challenges : MF vs. RPA or LDA vs. QSGW



Schilfgaarde, Kotani, Faleev PRL 96, 26402 (2006)

LDA is a mean field type of approach

$$\left( -\frac{\nabla^2}{2m} + V_{nuc}(r) + V_H(r) + V_{xc}^{LDA}(\rho(r)) \right) \Psi_{\mathbf{k}n}(r) = \epsilon_{\mathbf{k}n} \Psi_{\mathbf{k}n}(r)$$

GW approach is based on many-body theory

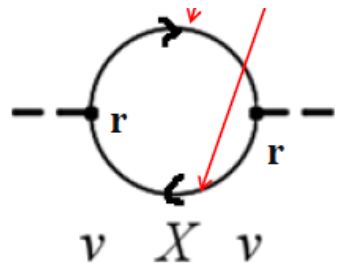
Typical error:

LDA 1-2 eV

G<sub>0</sub>W<sub>0</sub> 0.5 eV

QSGW 0.1-0.2 eV

$$\left( -\frac{\nabla^2}{2m} + V_{nuc}(r) + V_H(r) \right) \Psi_{\mathbf{k}n}(r) + \int dr \Sigma(r, r', \epsilon_{\mathbf{k}n}) \Psi_{\mathbf{k}n}(r') = \epsilon_{\mathbf{k}n} \Psi_{\mathbf{k}n}(r)$$



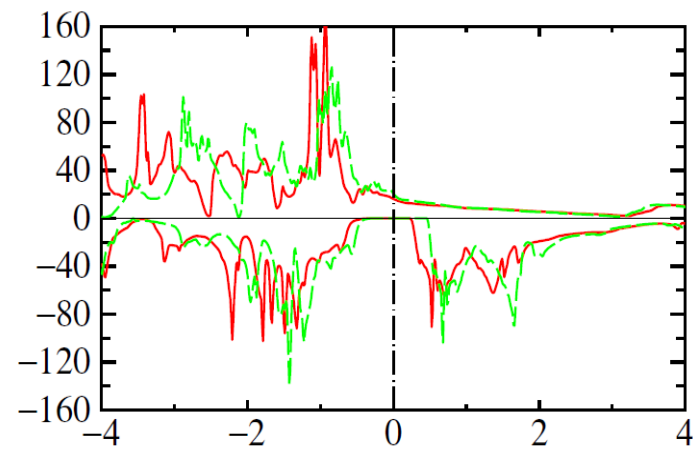
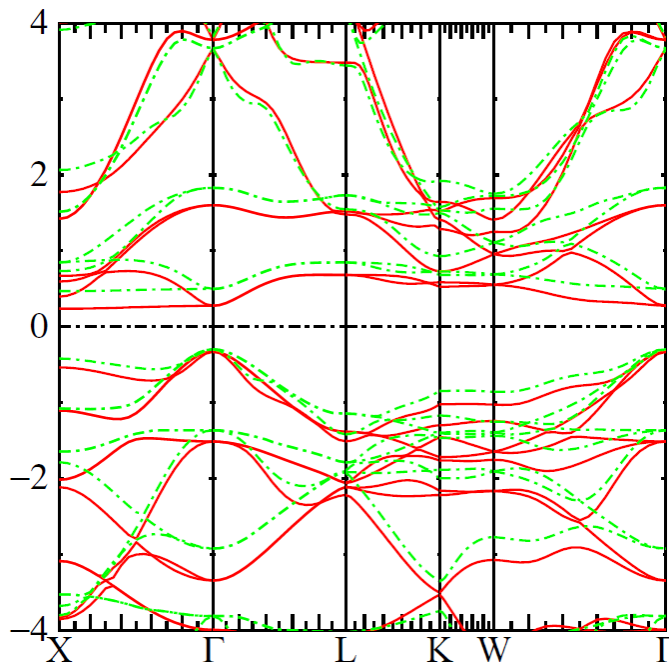
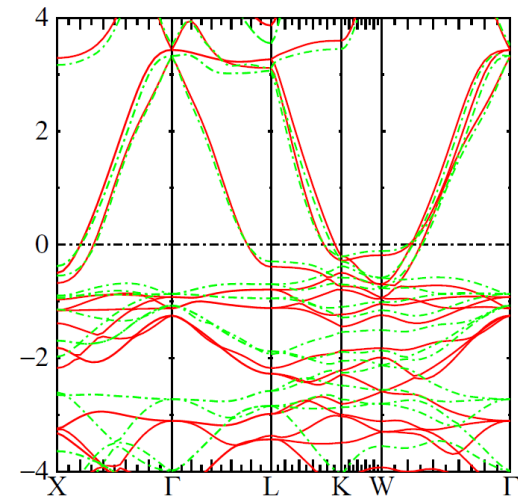
Polarizability (screening) !!!! RPA

# Systematic band "gap" analysis

Co<sub>2</sub>FeSi, Co<sub>2</sub>FeGe, Co<sub>2</sub>FeGa, Co<sub>2</sub>FeAl  
Co<sub>2</sub>MnSi, Co<sub>2</sub>MnGe, Co<sub>2</sub>MnGa, Co<sub>2</sub>MnAl

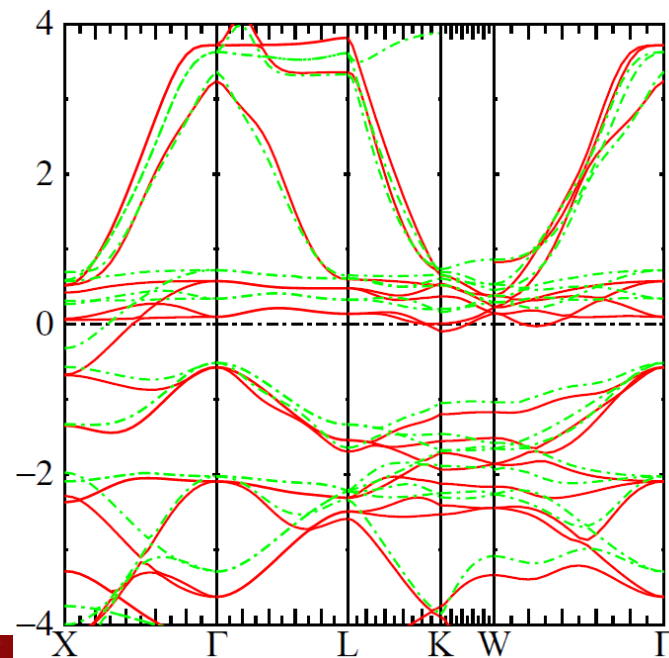
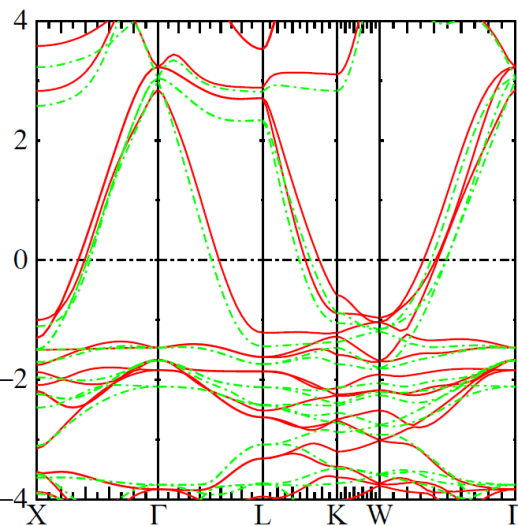
# Electronic structure : minority band gap

$\text{Co}_2(\text{Mn})\text{Si}$

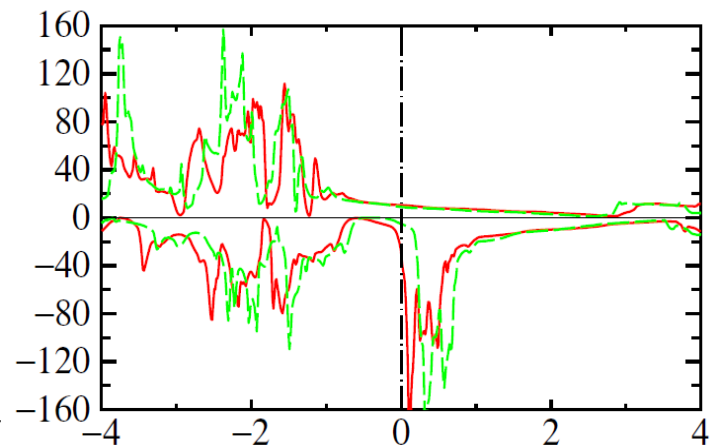


Co2MnSi : -LDA : 0.55 eV  
-GW : 0.7 eV

$\text{Co}_2(\text{Fe})\text{Si}$

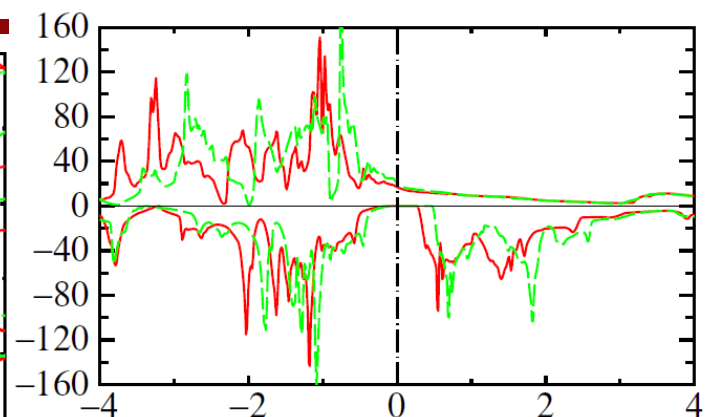
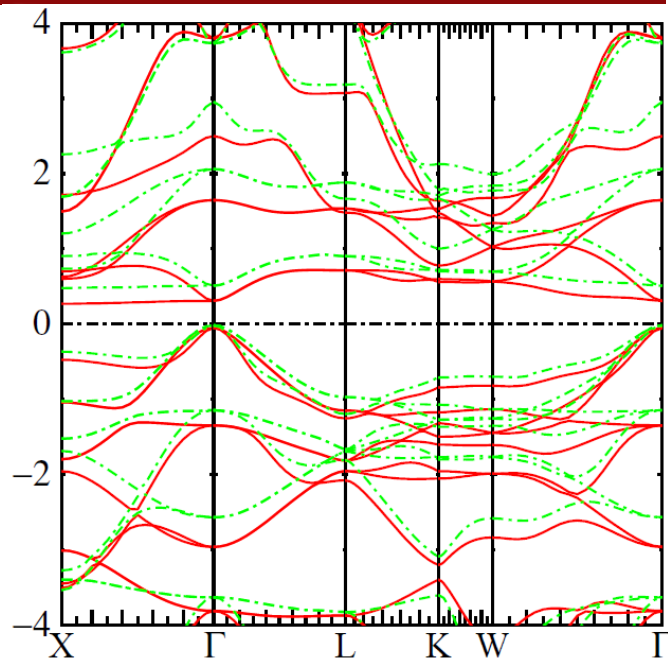
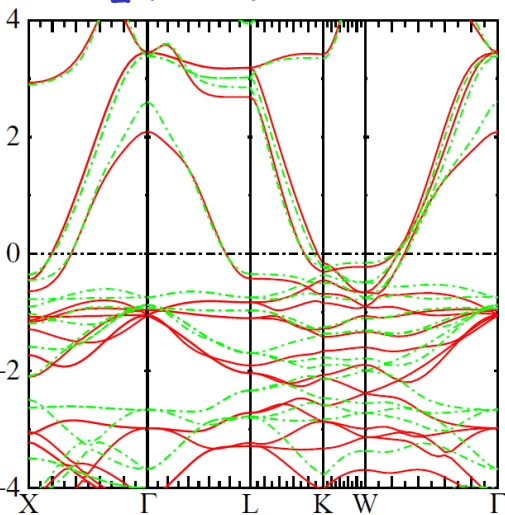


Co2FeSi : -LDA : 0.7 eV  
-GW : 1.33 eV



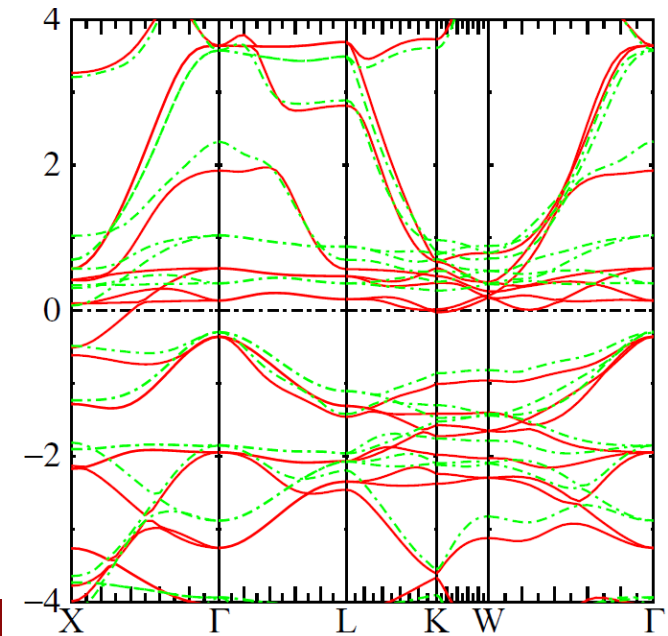
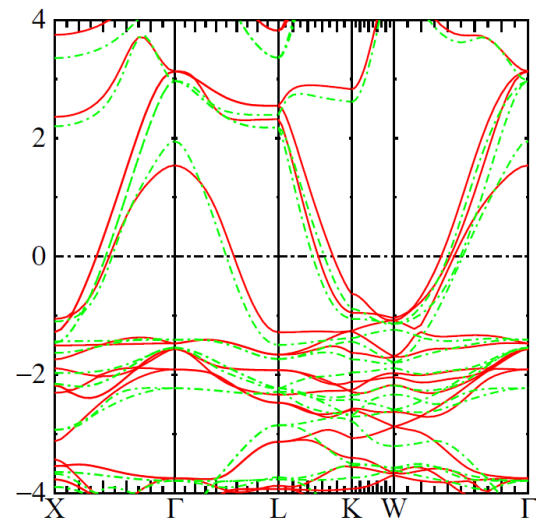
# Electronic structure : minority band gap

$\text{Co}_2(\text{Mn})\text{Ge}$

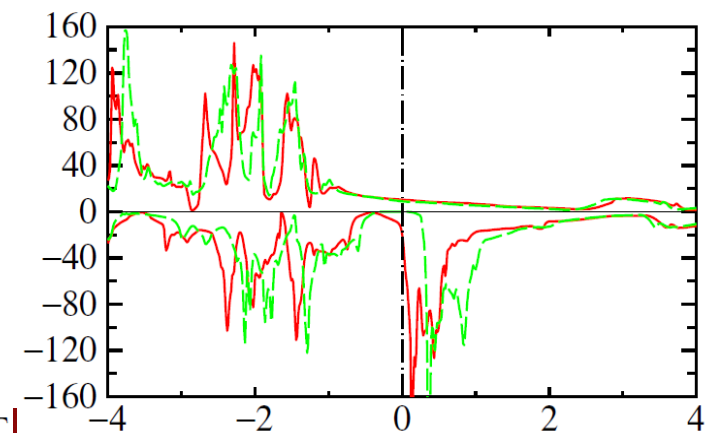


$\text{Co}_2\text{MnGe}$  : -LDA : 0.3 eV  
-GW : 0.45 eV

$\text{Co}_2(\text{Fe})\text{Ge}$

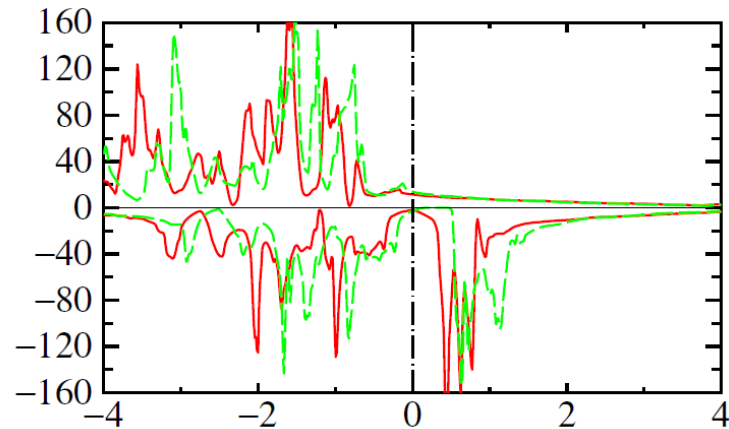
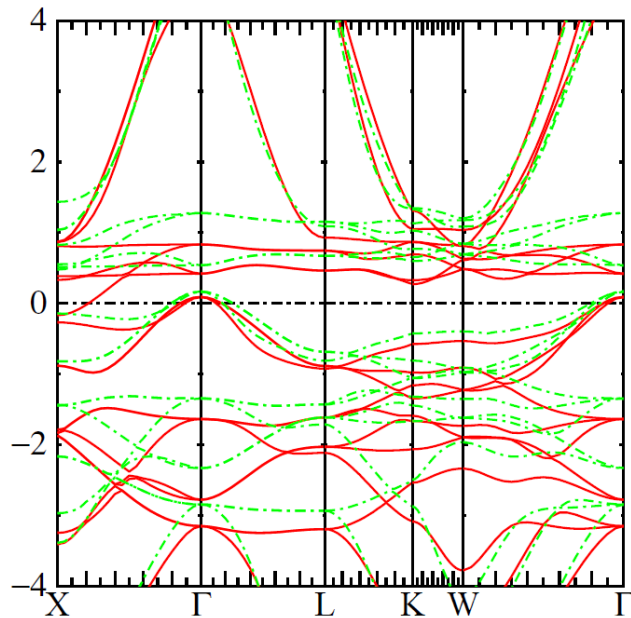
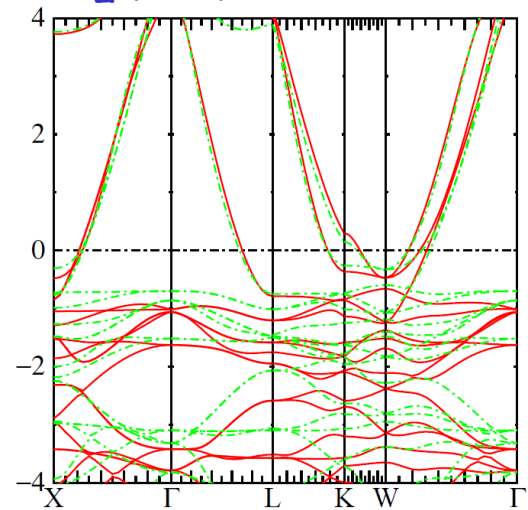


$\text{Co}_2\text{FeGe}$  : -LDA : 0.4 eV  
-GW : 0.6 eV



# Electronic structure : minority band gap

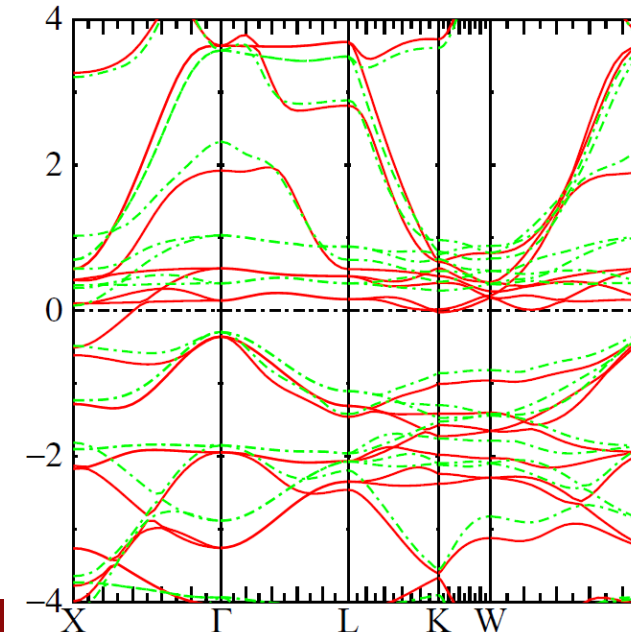
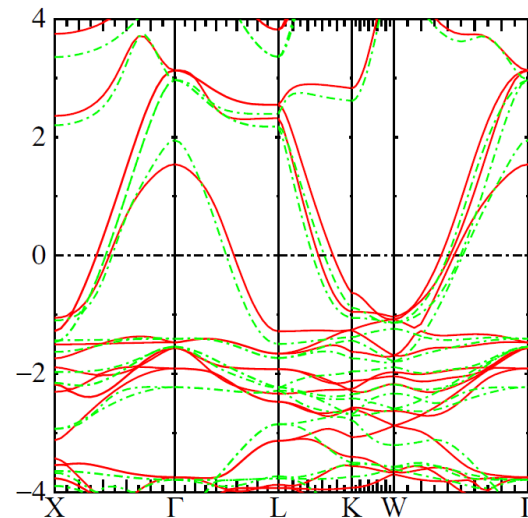
$\text{Co}_2(\text{Fe})\text{Ga}$



$\text{Co}_2\text{FeGa}$  : -LDA : 0.3 eV

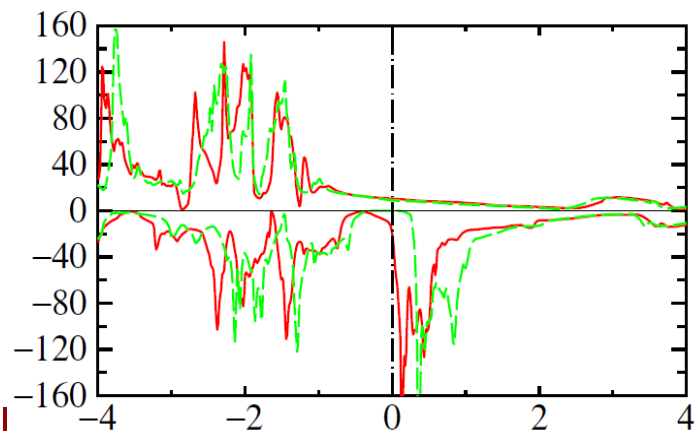
-GW : 0.3 eV

$\text{Co}_2(\text{Fe})\text{Ge}$



$\text{Co}_2\text{FeGe}$  : -LDA : 0.4 eV

-GW : 0.6 eV



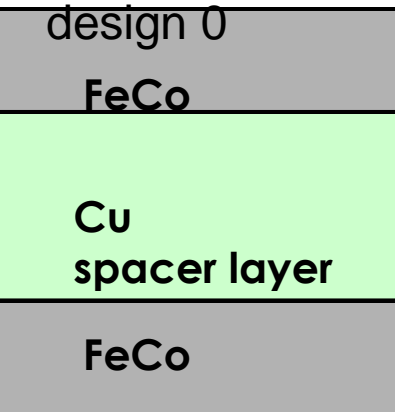


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- **K1 : Hexagonal Heusler like alloys :**
  - Magnetic Anisotropy
- **Summary and Conclusions**

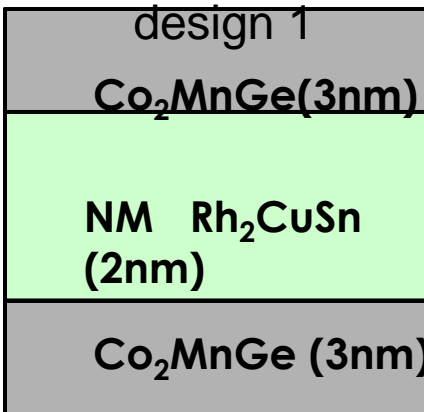
# GMR SVs with Heusler alloys

K. Nikolaev, P. Kolbo, T. Pokhil, X. Peng, Y. Chen, T. Ambrose, and O. Mryasov, *Appl. Phys. Lett.* **94**, (09);



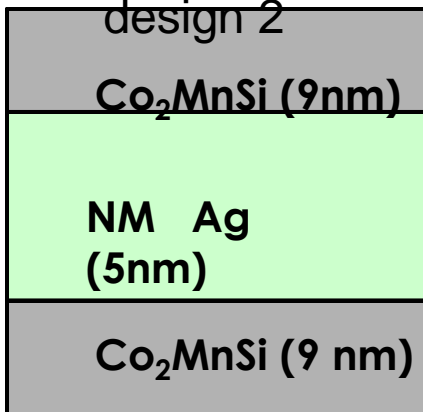
(111)

- **Co/Cu GMR (111) texture**
- Low-bias MR~3.5%;
- RA= 45mΩ-μm<sup>2</sup>
- ΔRA~ 1.6mΩ-μm<sup>2</sup>



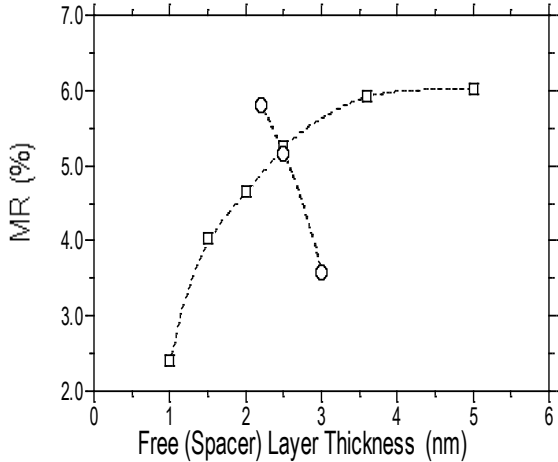
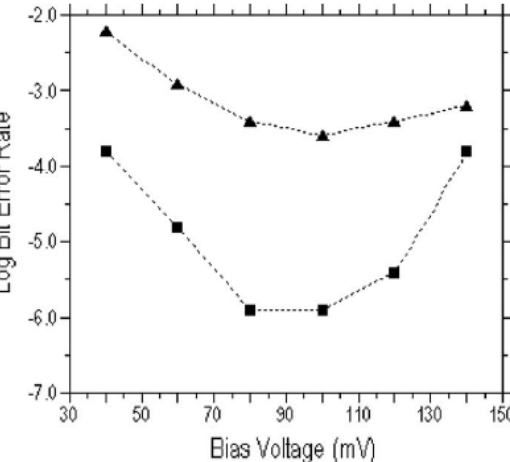
(110)

- **Ambrose, Mryasov US 6,876, 522 B2**
- **Nikolaev et.al. 2009, Seagate**
- Low-bias MR~6.8%;
- RA= 60mΩ-μm<sup>2</sup>
- ΔRA~ 4.0mΩ-μm<sup>2</sup>



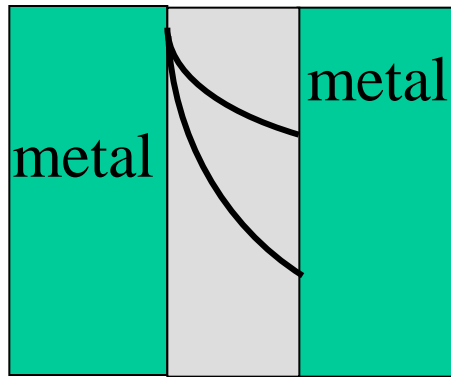
(100)

- **Iwase et.al. 2009,**
- Low-bias MR~28.8%;
- RA= 66mΩ-μm<sup>2</sup>
- ΔRA~8.9mΩ-μm<sup>2</sup>

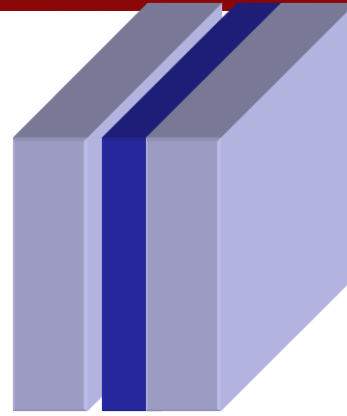


- **Design 1:**
  - anneal 200 C
  - test read heads
  - spacer short MFP
- **Design 2:**
  - problematic (100) texture
  - oxidation of spacer (Ag)
  - anneal seed 700 C/ CMS 350 C

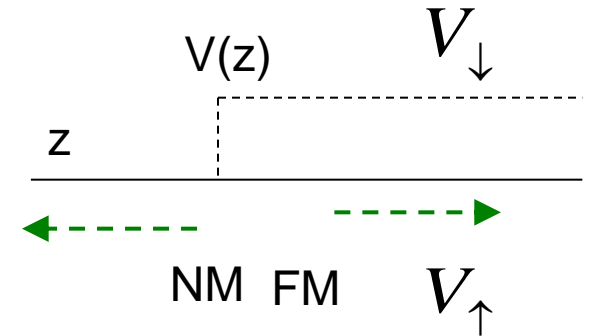
# BACKGROUND : spin dependent transport



TMR



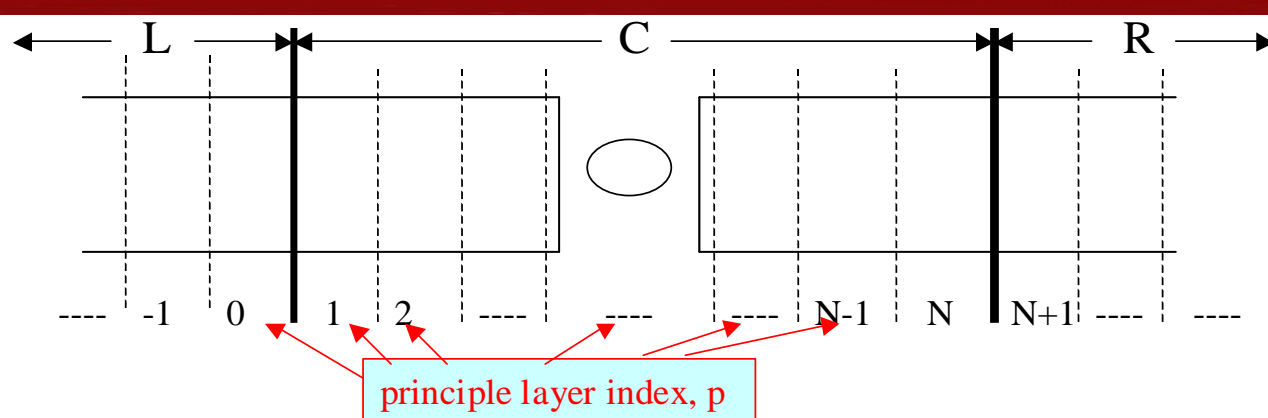
F/NM/F



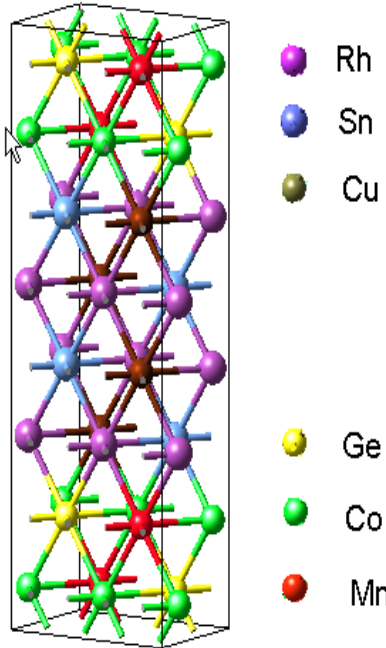
CPP-GMR

- if Modify materials set what RA and MR trends to expect
- Challenge is to improve spin dependent scattering :
  - $\beta$  and  $\gamma$  bulk and interface spin asymmetry
  - Curie point or spin mixing stability
- Compare with available experiment

# Direct transport simulations : model



$$\sigma \equiv \frac{dI}{dV} \Big|_{V=0} = -\frac{2e^2}{h} \text{Tr}[(\Sigma_{11} - \Sigma_{11}^+) g_{1N} (\Sigma_{NN} - \Sigma_{NN}^+) g_{N1}^+]_{E=E_F}$$



- All charge and potential relaxation takes place in central (C) region. Right (R) and left (L) regions assumed to have bulk charge density and potentials.
- Only adjacent layers interact, so the "Hamiltonian"  $h \equiv P(E) - S^\alpha$  of auxiliary Green's function  $g(E) = (P(E) - S^\alpha)^{-1}$  has nonzero matrix elements only for adjacent principle layers and takes tridiagonal form in principal layer index:

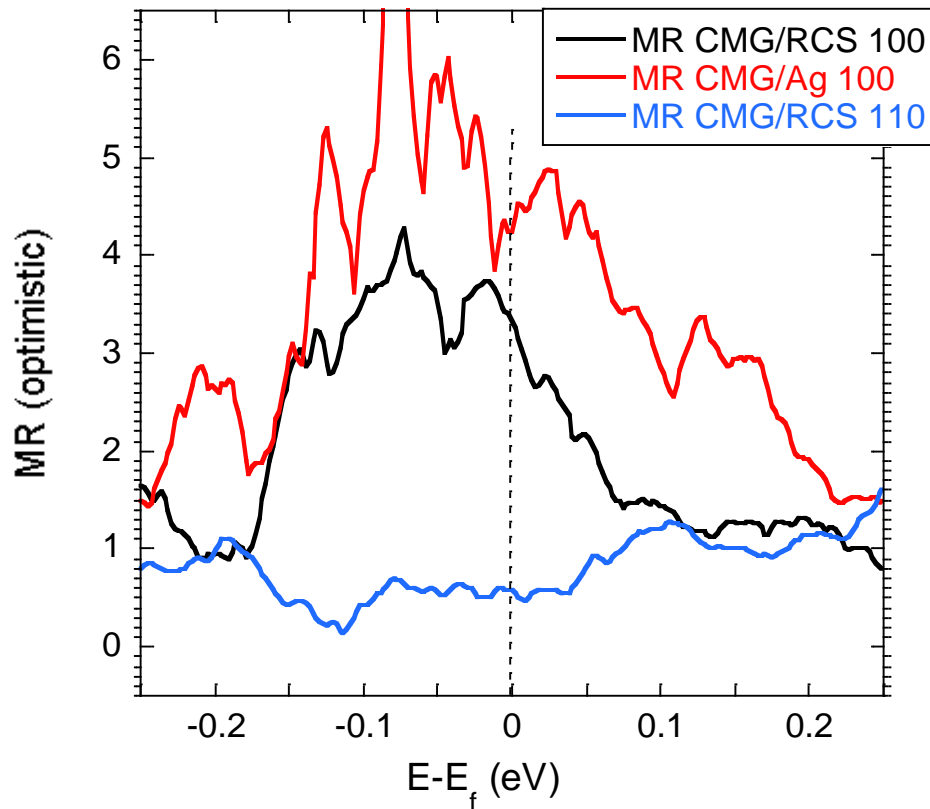
Model 1: L = Ag ; C = CMG|RCS|CMG ; R = Ag (110) and (100)  
 Ag|Co2MnGe(3)|Rh2CuSn(3)|Co2MnGe(3)|Ag

Model 2: L = Ag ; C = CMG|Ag|CMG ; R = Ag (100)  
 Ag|Co2MnGe(3)|Ag|Co2MnGe(3)|Ag

# Direct transport simulations : Design 1 vs. Design 2

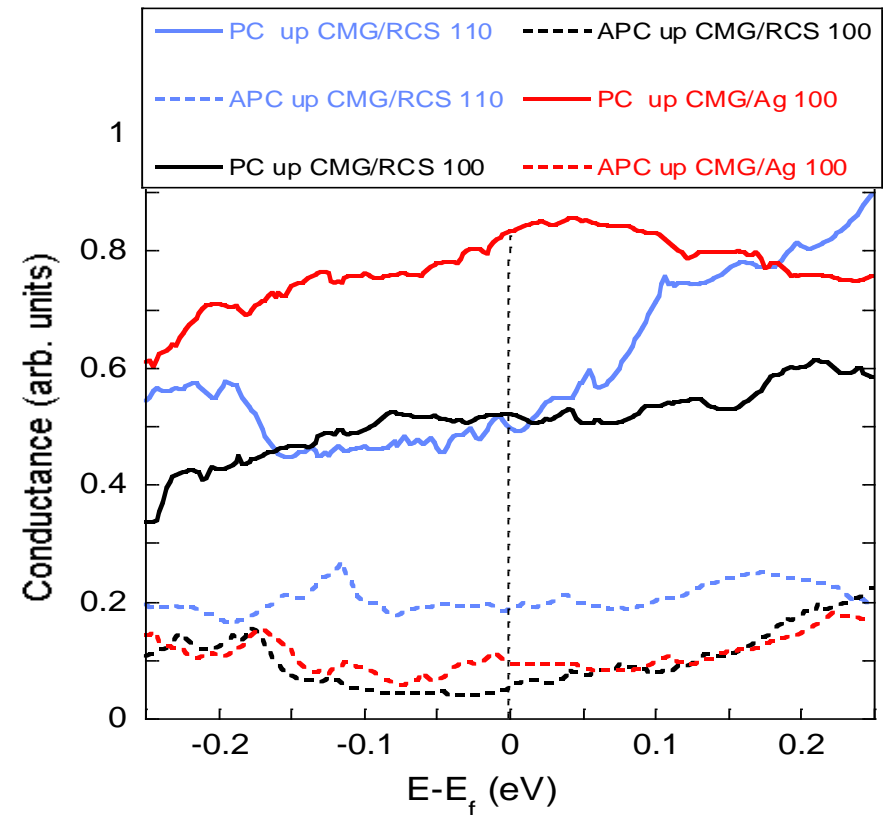
Model 1:

Ag|CMG(3)|RCS(3)|CMG(3)|Ag



Model 2:

Ag|CMG(3)|Ag|CMG(3)|Ag

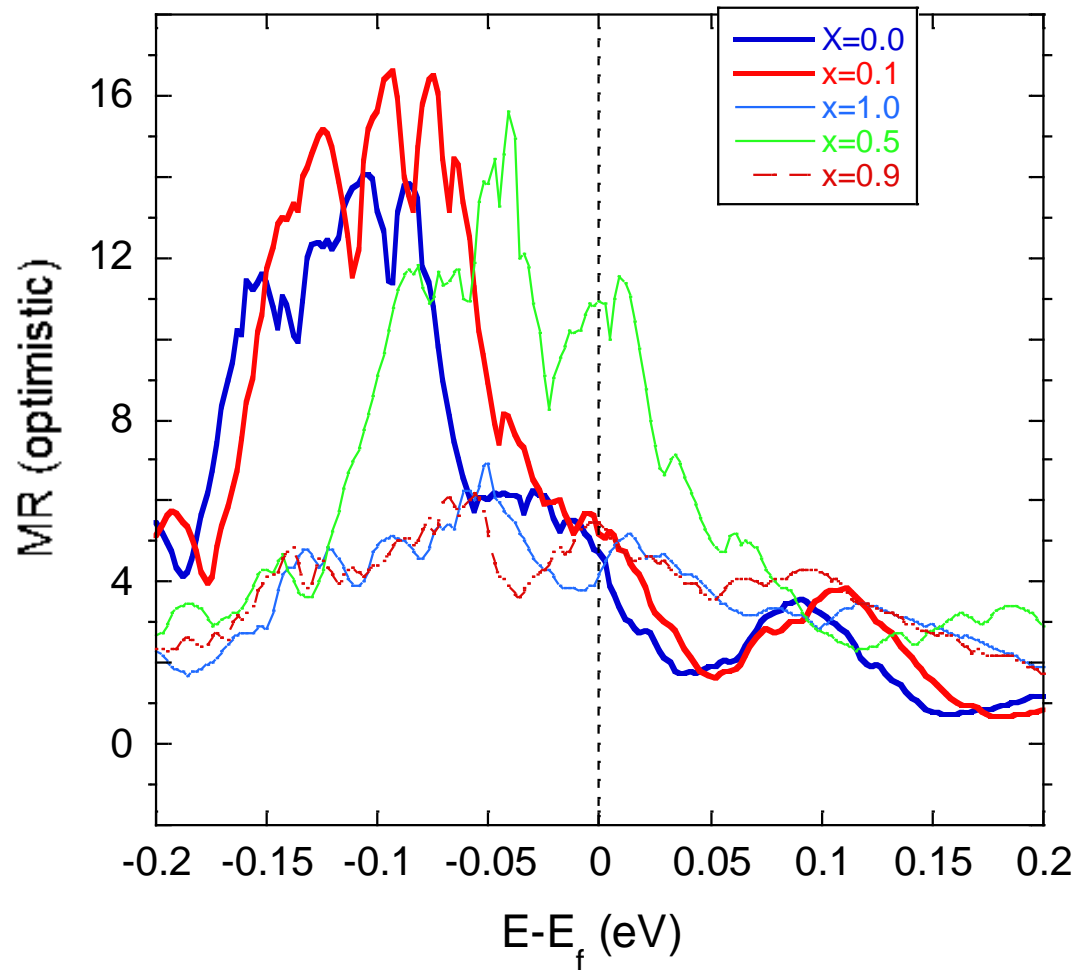


## Direct transport simulations findings:

- better conductance for majority in the case of CMG/Ag spacer than CMR/RCS
- > than 4x higher conductance in the minority channel for 110 texture than in 100

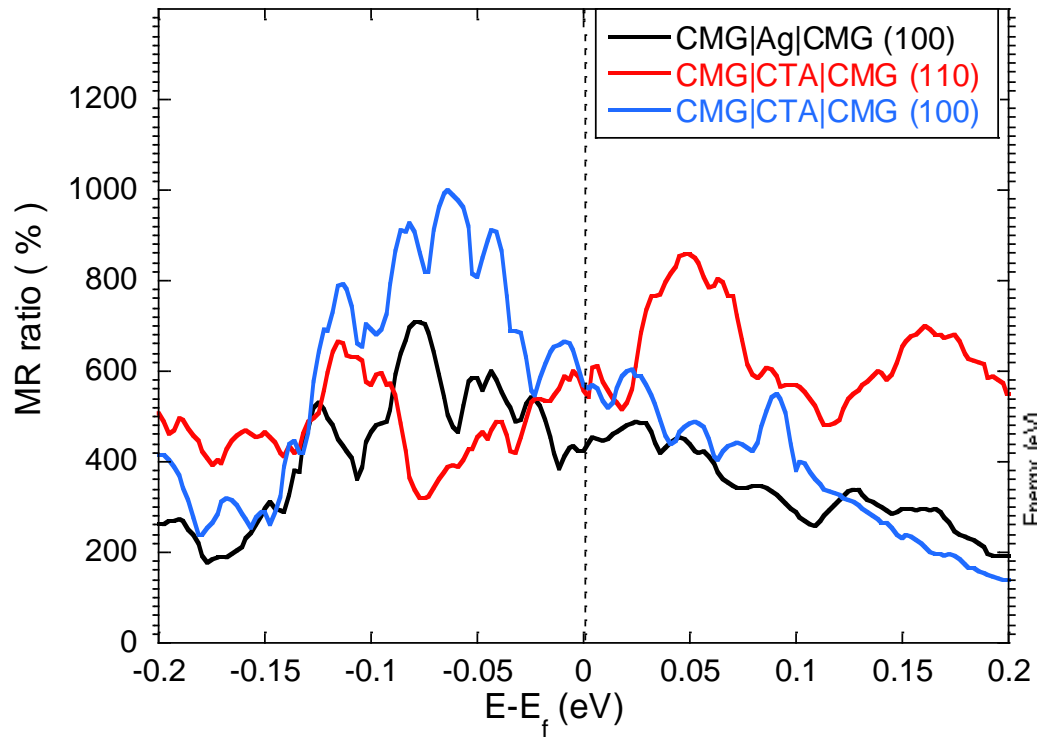
# Full transport simulations: $\text{Co}_2(\text{Mn-Fe})\text{Ge}$

## $\text{Co}_2(\text{Fe-Mn})\text{Ge}/\text{Ag}/\text{Co}_2(\text{F-M})\text{Ge}$ (001)

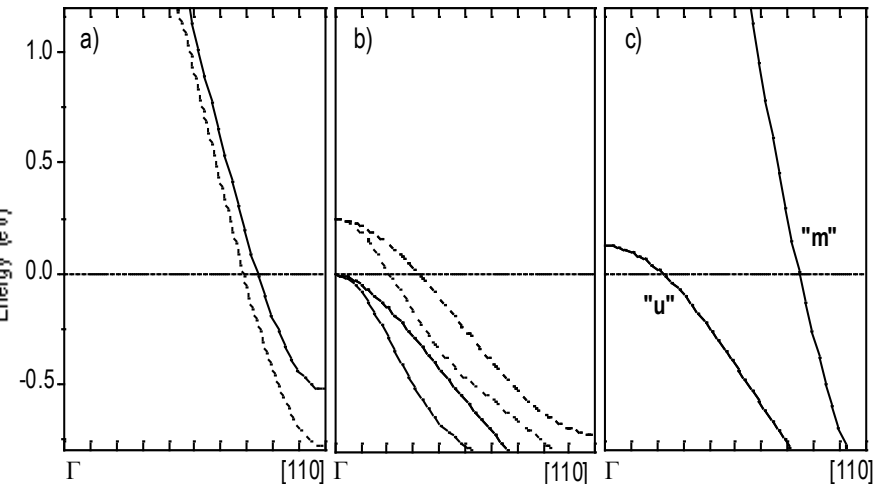


- $\text{Co}_2(\text{Fe-Mn})\text{Ge}$
- Preferable gaps
- Preferable  $T_c$
- Transport result support advantages of
- $\text{Co}_2(\text{Fe-Mn})\text{Ge}$

# Other Heusler alloys for all Heusler CPP-GMR



K. Nikolaev, P. Kolbo, T. Pokhil, X. Peng, Y. Chen, T. Ambrose, and O. Mryasov,  
*Appl. Phys. Lett.* **94**, (09);



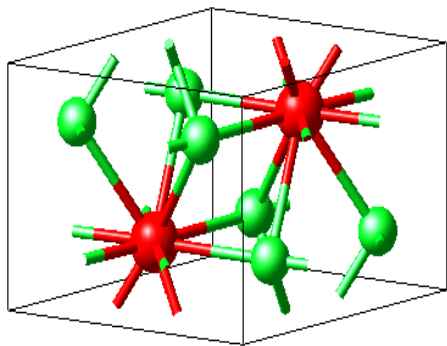
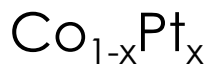
- within the ballistic transport limit all-Heusler junctions may outperform Ag based junction limited to (100) texture readily support more practical (110)
- experimental result in Prof. Hono's talk

# OUTLINE

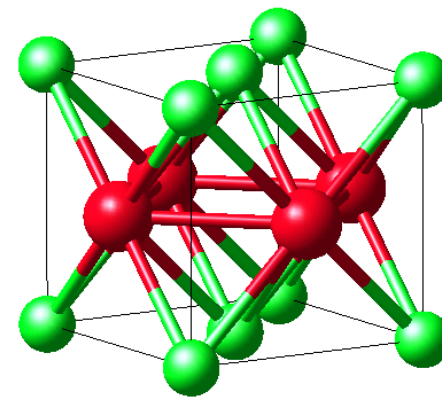
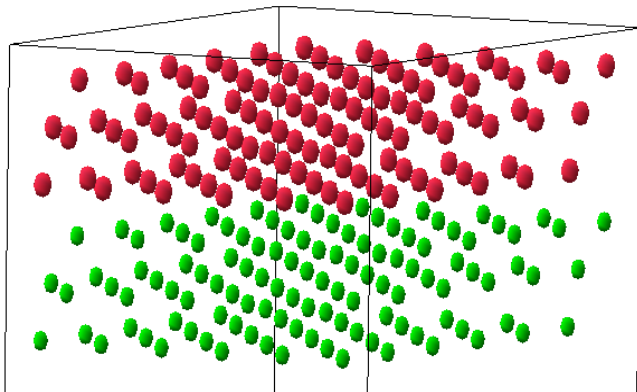
- $T_c$  : Comments on Disorder : spin disorder
  - spin mixing
  - Curie point calculations
  - quaternary alloys strategy
- $E_g$  : Electronic structure: minority band gap
  - fundamental gap theory
  - alloys design
- $\rho_{up}/\rho_{dn}$  : Spin Dependent Transport : GMR
  - Band matching
  - Q-alloy effects
- **K1** : Hexagonal Heusler like alloys :
  - Magnetic Anisotropy
- Summary and Conclusions



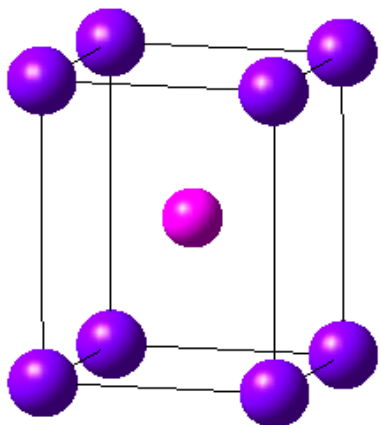
# Materials landscape : 3d-5d vs. 3d-metalloid



DO19



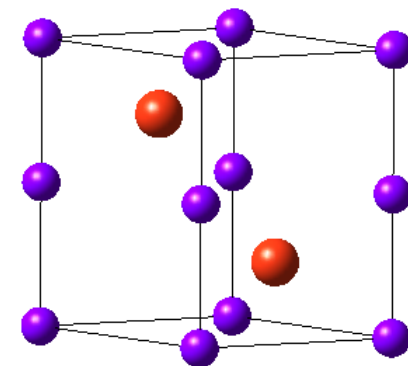
L10



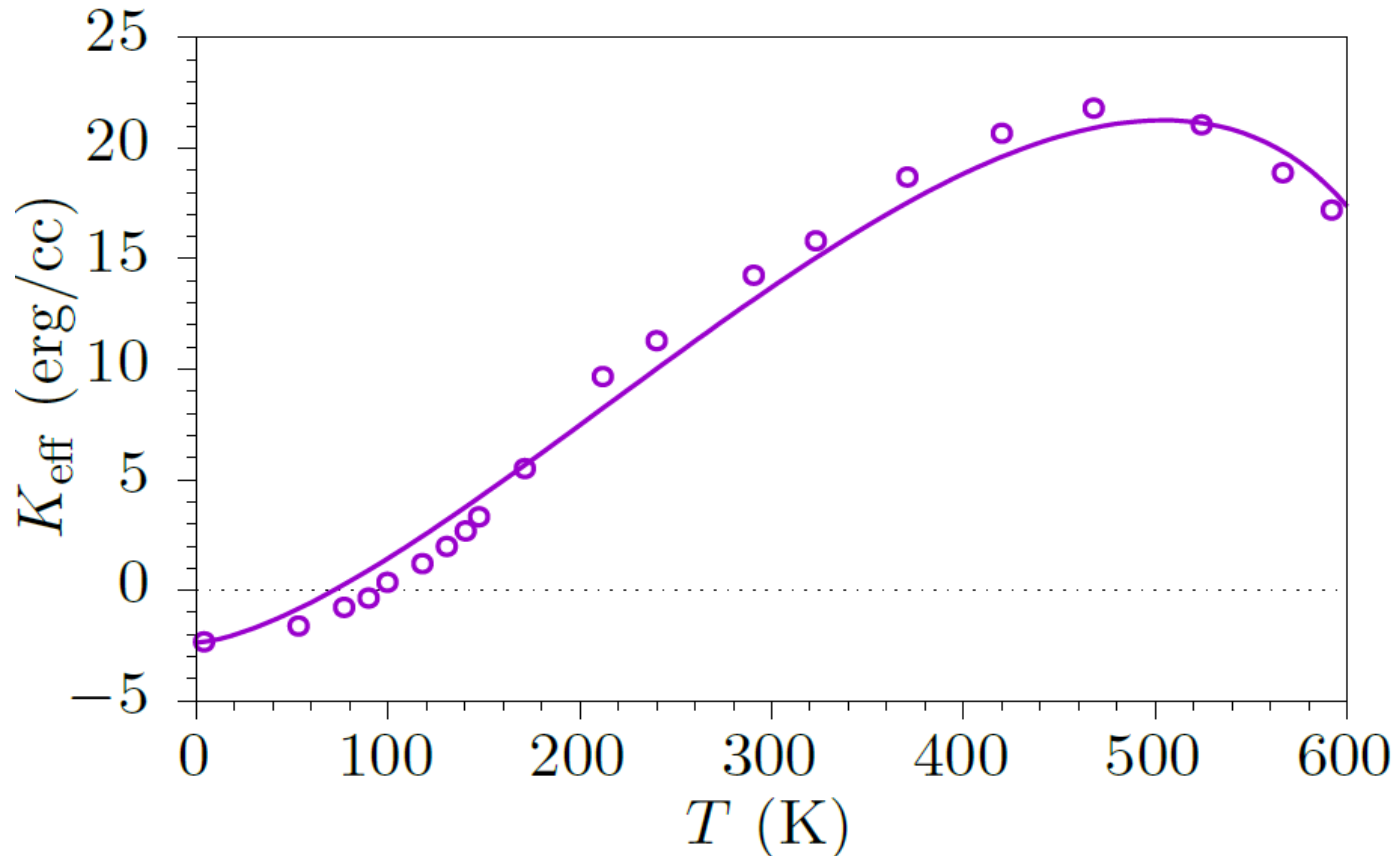
					B	C	N			
					Al	Si	P			
V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As
Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb
Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi
Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm

Two types of anisotropy:

- Bulk (FePt, Fe<sub>16</sub>N<sub>2</sub>, Mn-Bi)
- Interface (Fe/MgO)

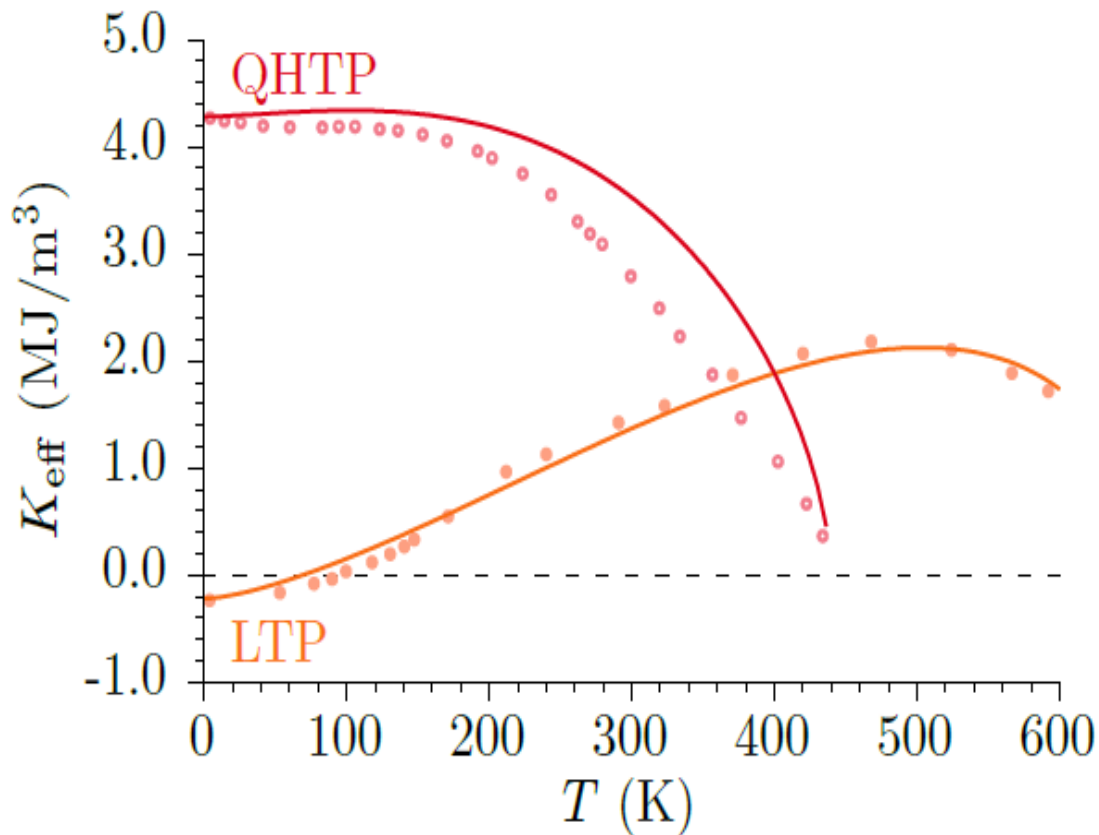


# $K(T)$ properties : the MnBi challenge



- Reorientation can be understood partially with lattice temperature
- Still hard to reconcile with large pick

# Our approach to the problem : $K_{\text{eff}}(T)$



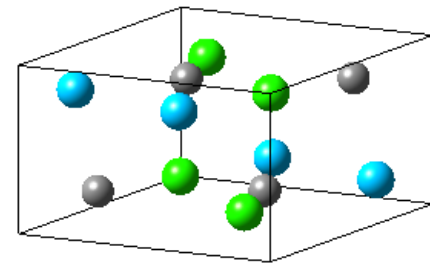
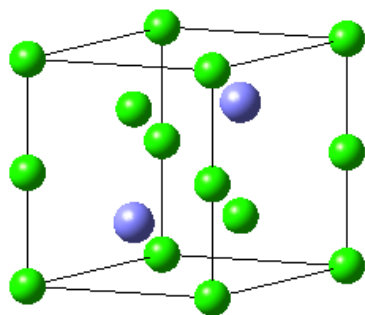
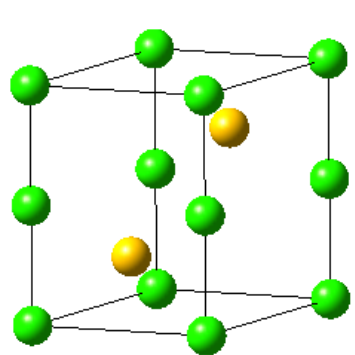
W. E. Stutius, T. Chen, and T. R. Sandin, AIP Conference Proceedings 18, 1222 (1974).

- $K_1(5\text{K}, K_2(5\text{K}), K_3(5\text{K})$  experimental
- determine  $k_2$  and  $d(2)$
- $m(T)$  from experiment

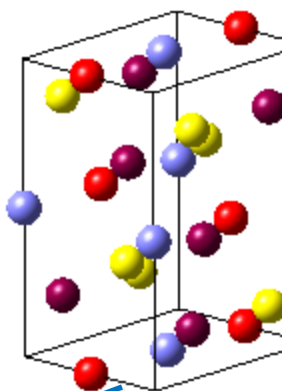
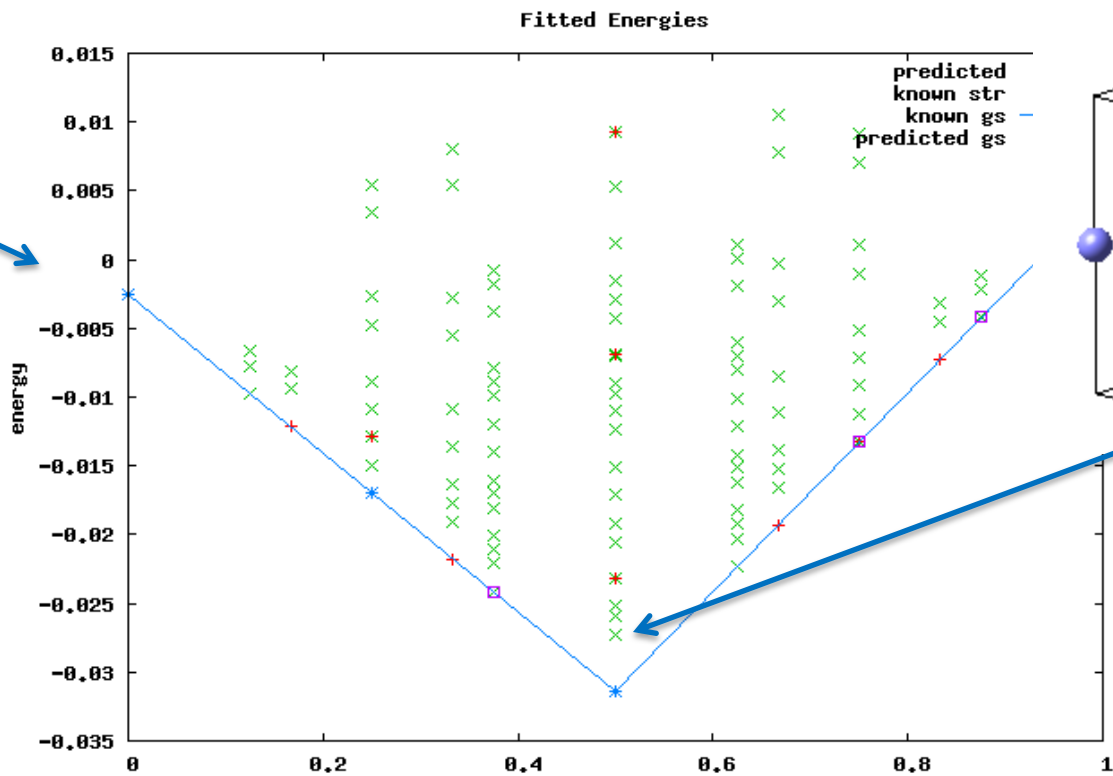
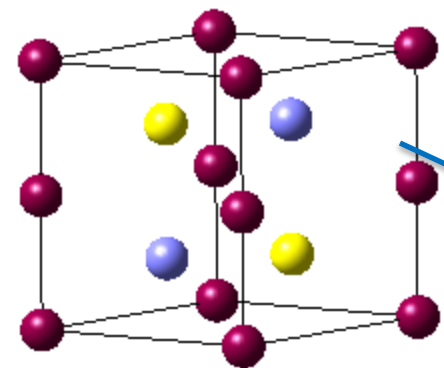
$$K_1 = \frac{3}{2}\kappa_2 + \delta_2 + \Delta_2$$

- Analysis indicate huge anisotropy of two-ion type  $d(2)$  of about + 8 MJ/mc compensated by large negative  $d(0)$  to give small in-plane  $K_{\text{eff}}$

# NiAs vs. Ni<sub>2</sub>In vs. TiNiSi : CE technique XYZ (Z=Ge)



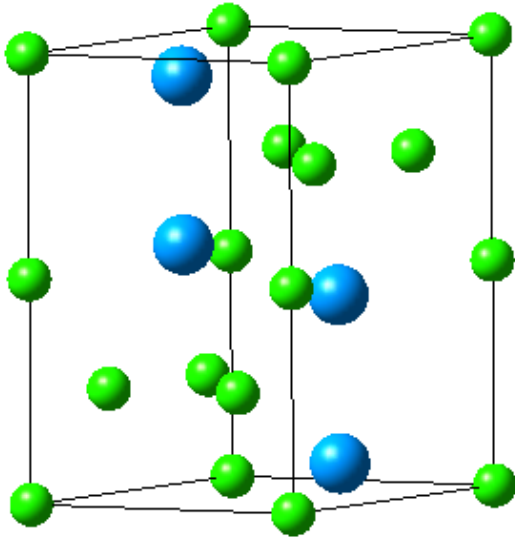
- Materials search : structure/phase generators: using generalized cluster expansion algorithms
- A. van de Walle, Nature Materials 7, 455 - 458 (2008)



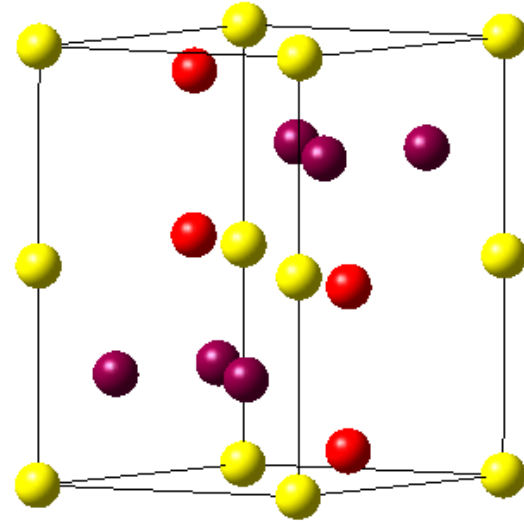
- Tuning Ms and Tc
- CoFe<sub>(1-x)</sub>Mn<sub>x</sub>Ge
- Too Low Tc
- K < 0

# NiAs vs. Ni<sub>2</sub>In vs. TiNiSi vs. Ni<sub>2</sub>U

Ni<sub>2</sub>U



XYZ , XY=MnFeCo, Z= Ge



- $K_{eff} = +0.45 \text{ MJ/m}^3$
- $M_s = 1310 \text{ emu/cc}$
- $T_c = 820 \text{ K}$

Decent performance , further enhancement of  $K$  might be needed

# Summary and Conclusions

- Addressed some of the material physics challenges:
  - (I)  $T_c$ , spin mixing ; (II) Band gap beyond DFT ;
  - (III) spin dependent transport ;
  - (IV) MAE
- 3d-3d hybridization nature of minority gap need to be addressed by the way of going beyond DFT
- Exchange coupling - MAE affected
- QSGW bands structure calculations enable accurate evaluation of HM features
  - $\text{Co}_2\text{FeGe}$  true HF (GW not LDA) with 0.6 eV gap
  - $\text{Co}_2\text{FeSi}$  X (GW) with 1.33 eV (0.7 eV) LDA gap unlike
- $\text{Co}_2(\text{Fe-Mn})\text{Ge}$  alloy shows favorable spin transport trend
- GMR CMG/CTA(RCS)/CMG(100) and (110) evaluated

# Acknowledgements:

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Prof. K. Hono : CTA based all Heusler spacer

2. MINT Center , The University of Alabama



*Joseph Barker<sup>2</sup>  
now AIMR Tohoku  
Assistant Prof.*



*Sergey Faleev<sup>2</sup>  
now IBM Almaden, SJ, CA*

- ARPA-E ; DARPA-SRC- C-SPIN ;  
early stage MnBi

C-SPIN



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Émergence/Partenariat  
Stratégique



*Alan Kalitsov<sup>1</sup>  
now WD*

BACK UP SLIDES

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